



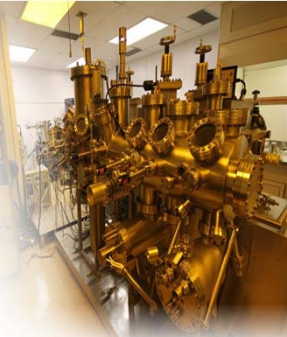
Confinement of Excitons in Strain-engineered Metamorphic InAs/In_xGa_{1-x}As/GaAs Quantum Dots for Long Wavelength (1.3 to 1.55 μm) Emission

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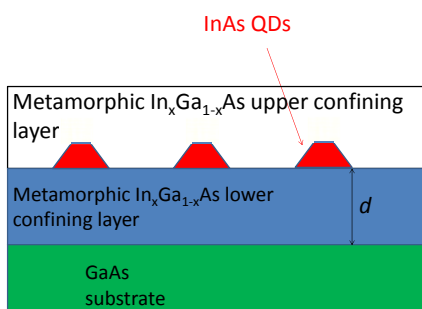


i. Introduction

Lasers with long wavelength emission (1.3 μm to 1.55 μm) are useful for telecommunications due to low attenuation in optical fibers. Although lasers based on self-assembled quantum dots (QDs) emitting at 1.3 μm are commercially available, producing longer wavelength devices has proven challenging. One possible route for achieving this is the use of strain-engineered InAs/In_xGa_{1-x}As/GaAs quantum dots grown on a metamorphic buffer of InGaAs, where In_xGa_{1-x}As is used as upper and lower confining layers for exciton confinement.

ii. How to Control the Long Wavelength Emission?

- (1) Changing the amount of indium "x" in the upper confining layer (UCL) and lower confining layer (LCL), which changes the band discontinuities between the QDs and confining layers and also affects the QD-CL mismatch (strain).
- (2) Changing thickness "d" of the LCL which affects only the mismatch between the QDs and LCL.¹⁻²



iii. Tool Used

We have studied thirty InAs/In_xGa_{1-x}As/GaAs QDs samples at 2 K and in magnetic fields up to 15 T. AFM measurements confirm that the dot morphology is independent of both *d* and *x*.

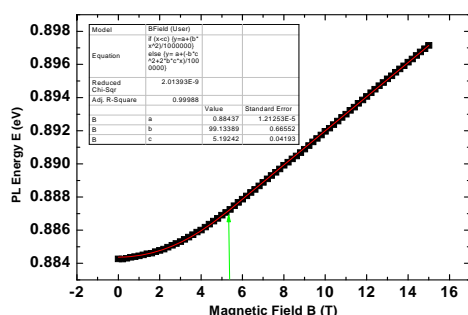
An excitonic model was used to calculate the properties of exciton, i.e., reduced mass, μ and bohr radius of exciton, a_B .

$$E = E_0 + e^2 a_B^2 B^2 / 8\mu \quad \text{for} \quad B \leq 2\hbar / ea_B^2 \quad \text{Eq-1}$$

$$E = E_0 - \hbar^2 / 2\mu a_B^2 + \hbar e B / 2\mu \quad \text{for} \quad B \geq 2\hbar / ea_B^2 \quad \text{Eq-2}$$

where a_B and μ are the Bohr radius and reduced mass of the exciton respectively.

The transition from Eq-1 to Eq-2 occurs when magnetic length becomes smaller than the spatial confinement.



iv. Our Data

Figure1: The circle size shows value of diamagnetic shift coefficient, $\Gamma = e^2 a_B^2 / 8\mu$. The blue color represents samples for which high field regime (where the magnetic length is $< a_B / \sqrt{2}$) is reached within available fields and the red color where it is not.

Figure 2: Bohr radius of exciton, a_B , as a function of QD-CL mismatch.

Figure 3: Reduced mass of exciton, μ , as a function of QD-CL mismatch.

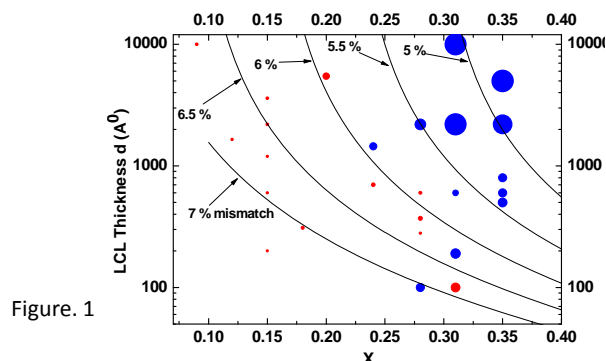


Figure. 1

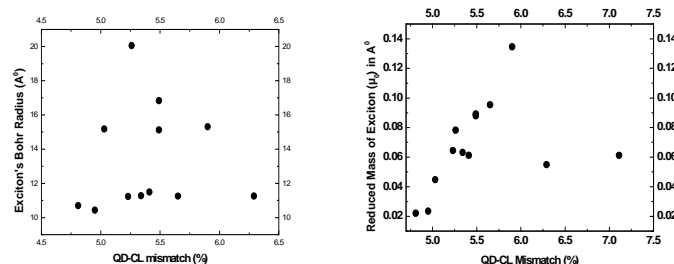


Figure. 2

Figure. 3

v. Observation, Discussion and Conclusion

We observe in Fig. 1 that for samples where *x* is low (≤ 0.24) the high-field regime is not reached irrespective of *d*, but does when it is high (≥ 0.24 %). Γ increases with increasing *x* when *d* is constant but shows complex behaviour with increasing *d* when *x* is constant. First it decreases with increasing *d*, then increases and saturates.

Fig. 3 shows that reduced mass of exciton, μ increases with increasing mismatch for samples where the mismatch is low (5 to 6 %) and then drops dramatically, but there is no clear trend in a_B with mismatch in Fig. 2.

Further Work:

To study more samples with parameters indicated by "☺" in Figure 4 to understand the trend in Γ at constant *x*.

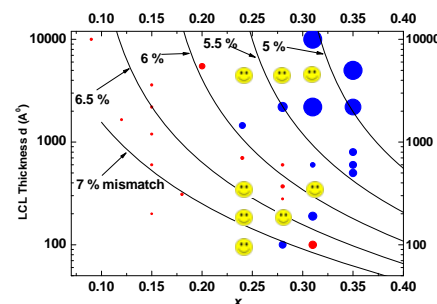


Figure:4