# Spin-photon quantum interface in quantum dots

A. Imamoglu

Quantum Photonics Group, Department of Physics

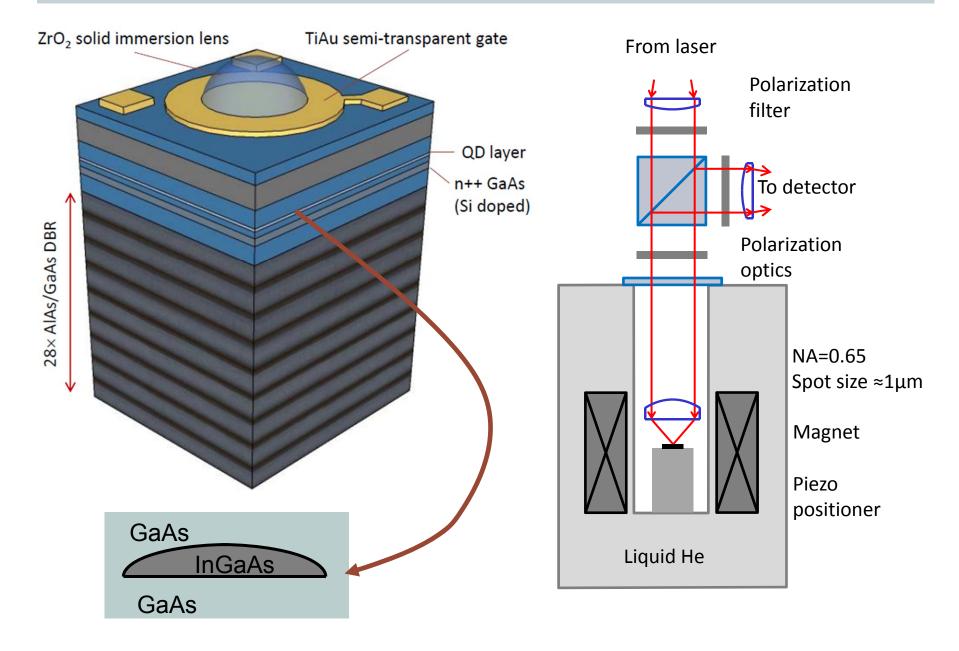
ETH-Zürich

- 1. Spin-photon entanglement (see also Peter McMahon's talk)
- 2. Quantum teleportation of a photonic qubit to a spin qubit

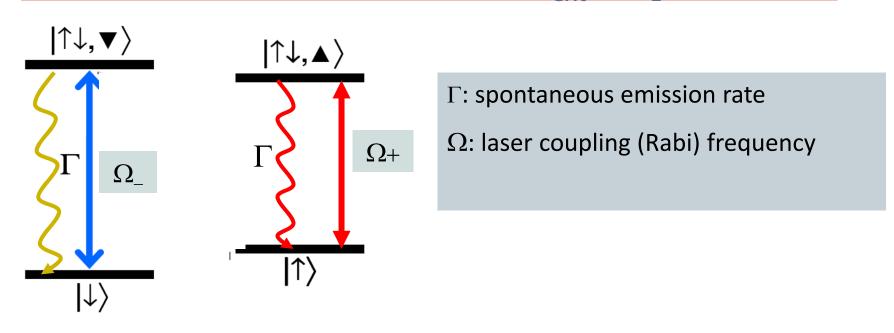
## Spin-photon quantum interface

- GaAs based semiconductors exhibit highly efficient spin-dependent optical transitions
- ⇒ optical manipulation of spin qubits
- Quantum dots in photonic nanostructures allow for efficient extraction of photons
- ⇒ high efficiency single-photon sources

## Resonant quantum dot Spectroscopy

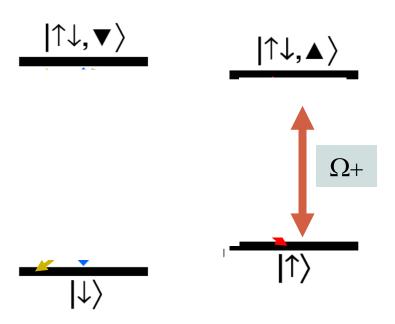


## Strong spin-polarization correlations: Faraday geometry ( $B_{ext} = B_z$ )



- QD with a spin-up (down) electron only absorbs and emits  $\sigma$ + ( $\sigma$ -) photons a recycling transition similar to that used in trapped ions.
  - ⇒ Spin measurement

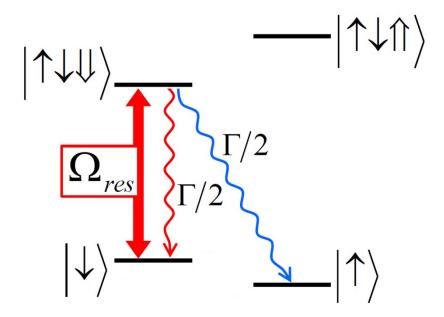
### Strong spin-polarization correlations: Faraday geometry ( $B_{ext} = B_7$ )



- QD with a spin-up (down) electron only absorbs and emits  $\sigma$ + ( $\sigma$ -) photons a recycling transition similar to that used in trapped ions.
  - ⇒ Spin measurement
- An off-resonant σ+ laser causes ac-Stark shift only for the |↑> state, acting as an effective magnetic field along the z-direction.

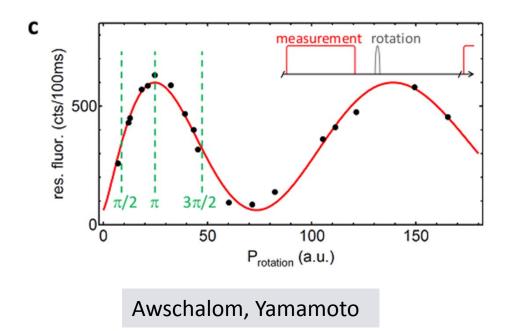
# Different selection rules in Voigt geometry $(B_{ext} = B_x)$

Excitation of a trion state results in either emission of a H polarized red photon to  $|\downarrow\rangle$  state or a V polarized blue photon to  $|\uparrow\rangle$  state, with equal probability.



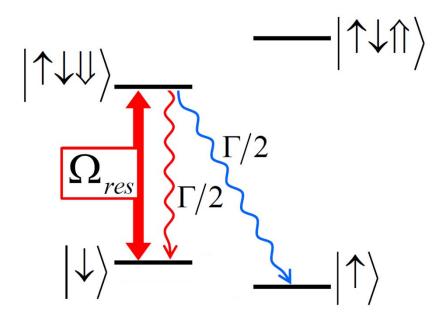
# Spin rotation using off-resonant circularly polarized lasers

• External field along x ( $B_{ext} = B_x$ ): quantization axis orthogonal to the laser-induced effective field



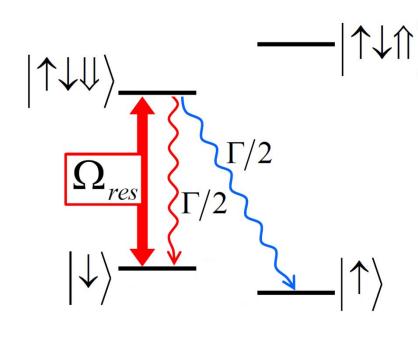
# Selection rules in Voigt geometry $(B_{ext} = B_x)$

Excitation of a trion state results in either emission of a H polarized red photon to  $|\downarrow\rangle$  state or a V polarized blue photon to  $|\uparrow\rangle$  state, with equal probability.



# Selection rules in Voigt geometry $(B_{ext} = B_{x})$

Excitation of a trion state results in either emission of a H polarized red photon to  $|\downarrow\rangle$  state or a V polarized blue photon to  $|\uparrow\rangle$  state, with equal probability.

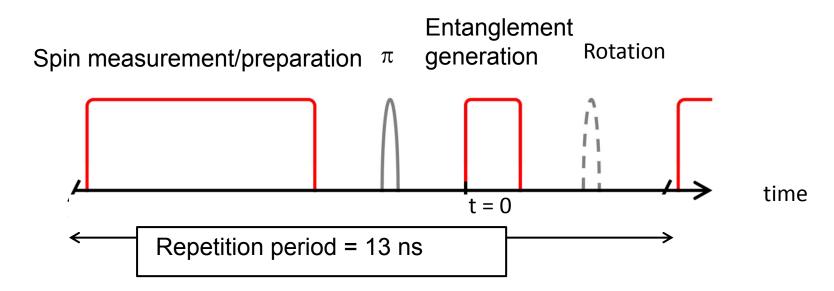


⇒ Spin-photon entanglement:
 potentially near-deterministic
 entanglement generation at
 ~1 GHz rate

$$|\Psi\rangle = \frac{1}{\sqrt{2}}(|\downarrow\rangle|\omega_{red};H\rangle + i|\uparrow\rangle|\omega_{blue};V\rangle)$$

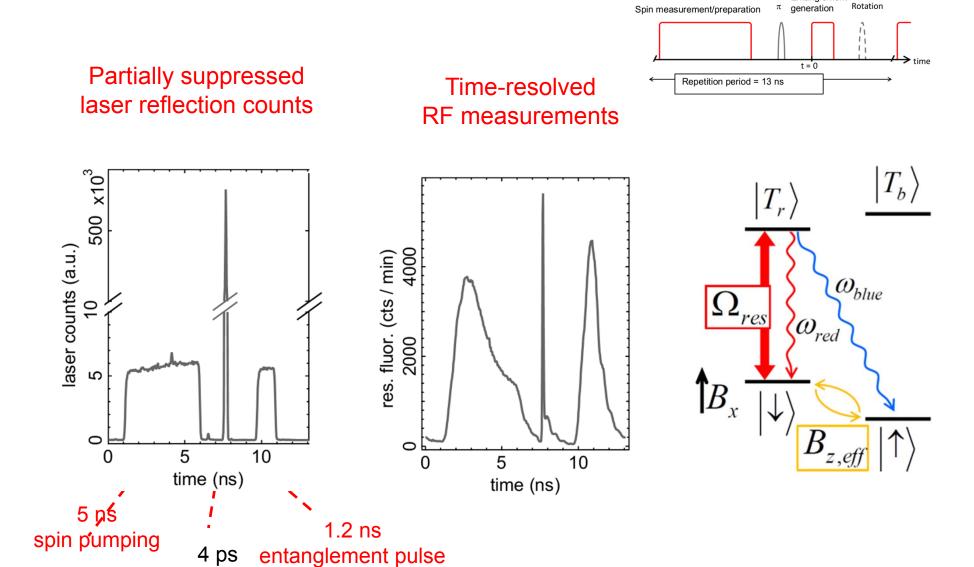
Similar results by Yamamoto, Steel groups; earlier work by Monroe, Lukin

# Procedure for spin-photon entanglement generation



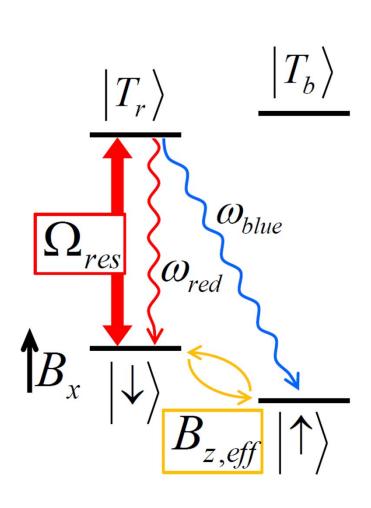
### Time resolved resonance fluorescence (RF)

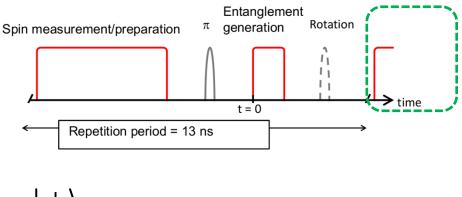
Entanglement



Rotation pulse

## Spin measurement and pumping

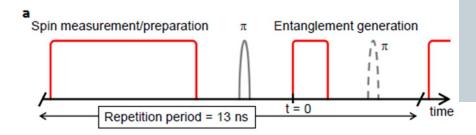




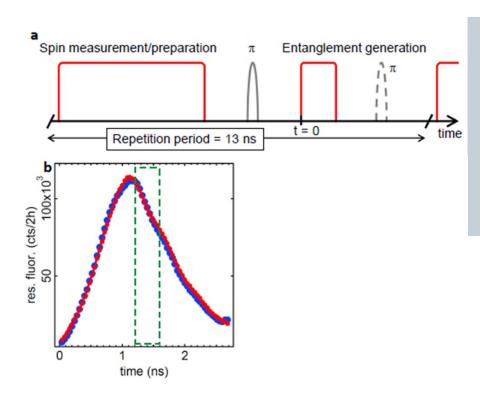
$$\left|\downarrow\right\rangle$$
  $\rightarrow$  ~2 photons/pulse.

$$|\uparrow\rangle$$
  $\rightarrow$  Nothing

- The detection of a photon shows the spin is in the state  $|\downarrow\rangle$
- At the end of the pulse, the spin is prepared in  $|\uparrow\rangle$

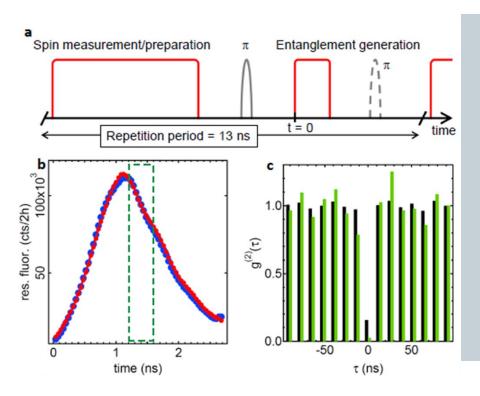


An additional  $\pi$ -pulse (dashed curve) is applied to realize a heralded measurement in the spin-up state.



An additional  $\pi$ -pulse (dashed curve) is applied to realize a heralded measurement in the spin-up state.

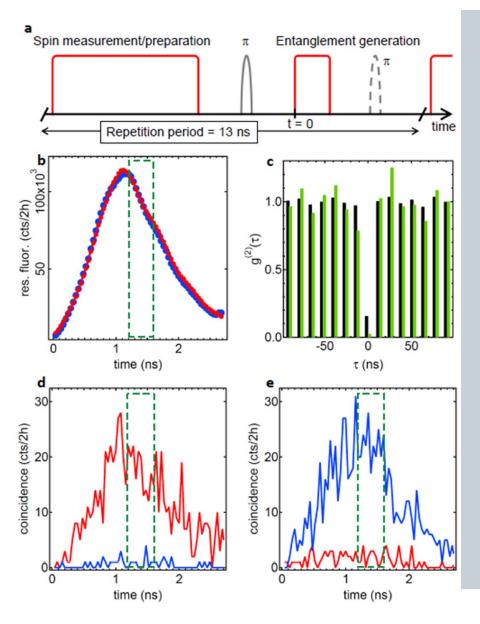
Identical (unconditional) counts for red and blue photons confirm the selection rules.



An additional  $\pi$ -pulse (dashed curve) is applied to realize a heralded measurement in the spin-up state.

Identical (unconditional) counts for red and blue photons confirm the selection rules.

The g(2) measurement shows that for the [1.2ns, 1.64ns] time range, probability of two-photon emission is negligible.



An additional  $\pi$ -pulse (dashed curve) is applied to realize a heralded measurement in the spin-up state.

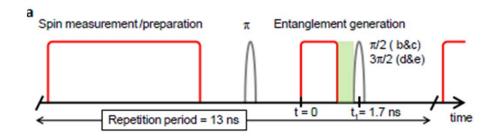
Identical (unconditional) counts for red and blue photons confirm the selection rules.

The g(2) measurement shows that for the [1.2ns, 1.64ns] time range, probability of two-photon emission is negligible.

A spin down (up) measurement event ensures that the detected photon is red (blue).

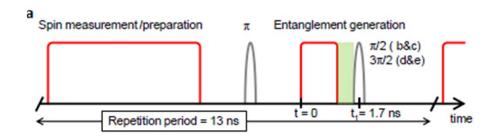
F1=0.87 ± 0.05 in the computational basis measure.ment

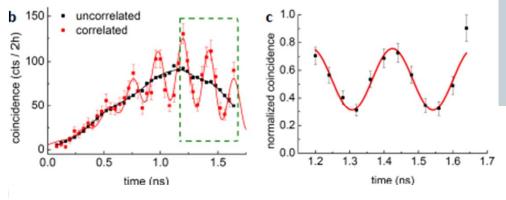
#### Measurement of quantum correlations



- An additional  $\pi/2$  or  $3\pi/2$ pulse (dashed curve) is
applied to measure the spin
in  $|\uparrow\rangle \pm |\downarrow\rangle$ .

#### Measurement of quantum correlations



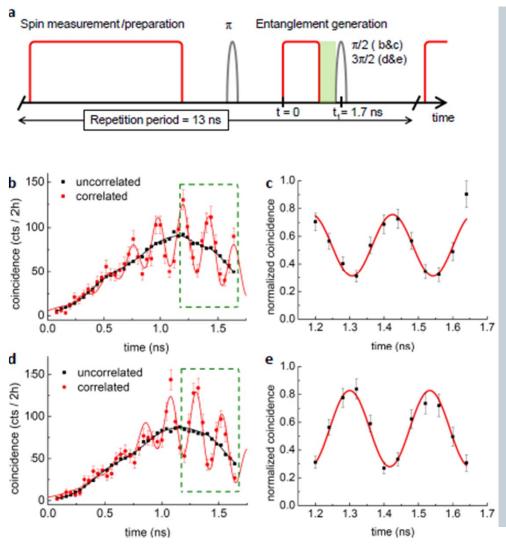


- An additional  $\pi/2$  or  $3\pi/2$ pulse (dashed curve) is
  applied to measure the spin
  in  $|\uparrow\rangle \pm |\downarrow\rangle$ .
- The data in b & c shows the coincidence measurement when  $\pi/2$ -pulse is applied.

$$|\tilde{\Phi}\rangle = \frac{1}{\sqrt{2}}(|\omega_{red}\rangle e^{-i\omega_z(t_1-t_g)} - i|\omega_{blue}\rangle)$$

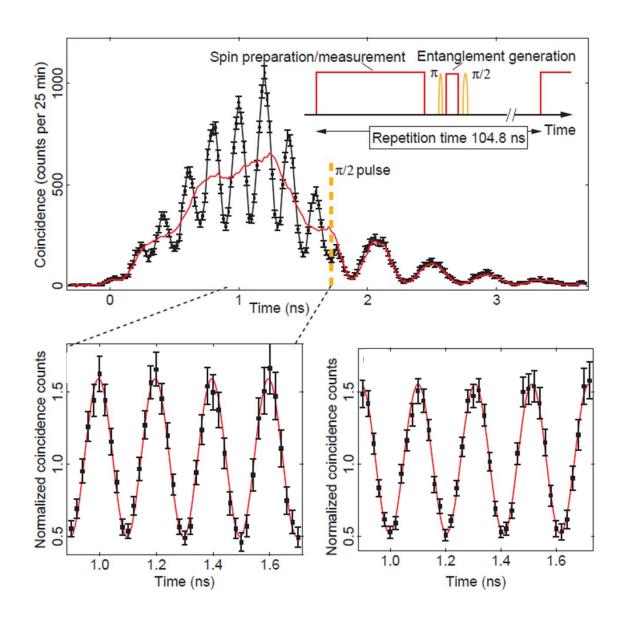
⇒ Photon generation events at different times correspond to a measurement of the photonic wave-function in different basis.

#### Measurement of quantum correlations



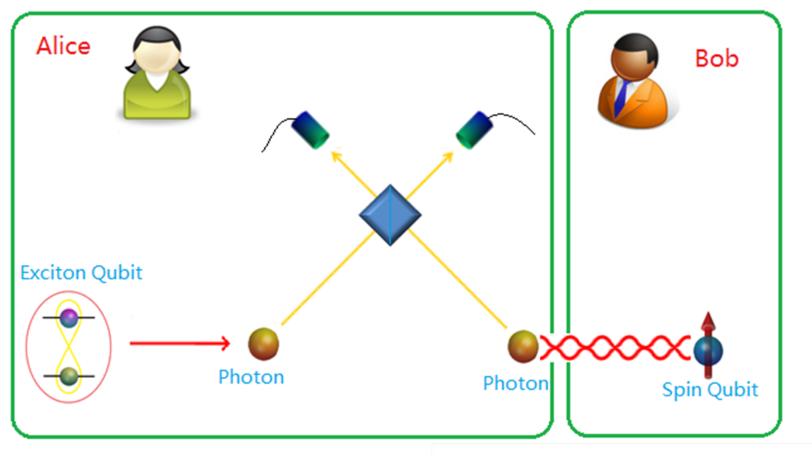
- An additional  $\pi/2$  or  $3\pi/2$ pulse (dashed curve) is
  applied to measure the spin
  in  $|\uparrow\rangle \pm |\downarrow\rangle$ .
- The data in b & c shows the coincidence measurement when  $\pi/2$ -pulse is applied.
- The data in d & e shows the coincidence measurement when  $3 \pi/2$ -pulse is applied.
- F2=0.46 ± 0.04 in the rotated basis measurement; overall fidelity F = 0.67 ± 0.05

## Improved spin-photon entanglement



- Fidelity limited primarily by the detector jitter
- The oscillation period before (after) the π/2 pulse is given by electron (hole)
   Zeeman energy.

# Teleportation from a photonic qubit to a solid-state spin qubit

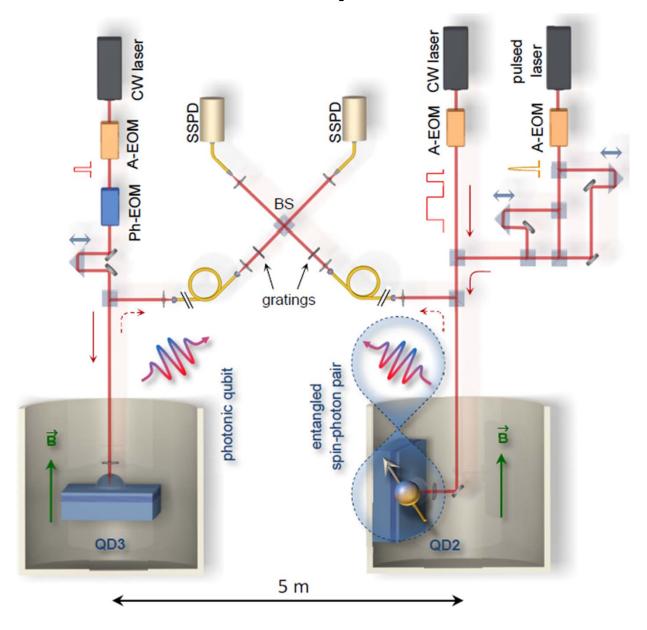


$$|\psi_p>=\alpha|\omega_b>_A+\beta|\omega_r>_A$$

$$|\Psi\rangle = \frac{1}{\sqrt{2}}(|\downarrow\rangle|\omega_{red};H\rangle + i|\uparrow\rangle|\omega_{blue};V\rangle)$$

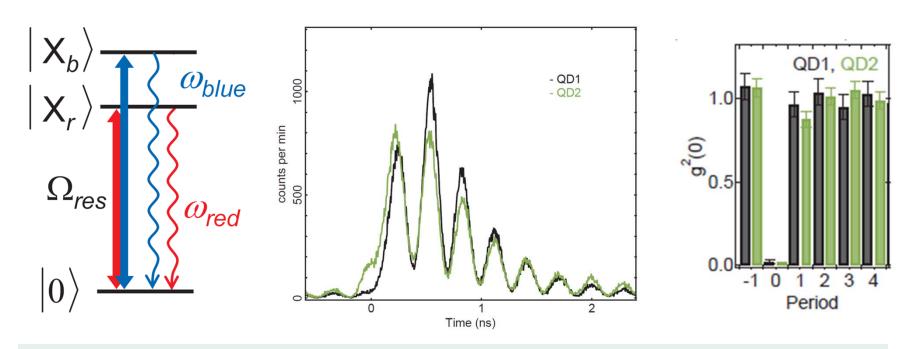
Output:  $|\psi_e>=\alpha|\downarrow>+\beta|\uparrow>$ 

## Experimental teleportation scheme



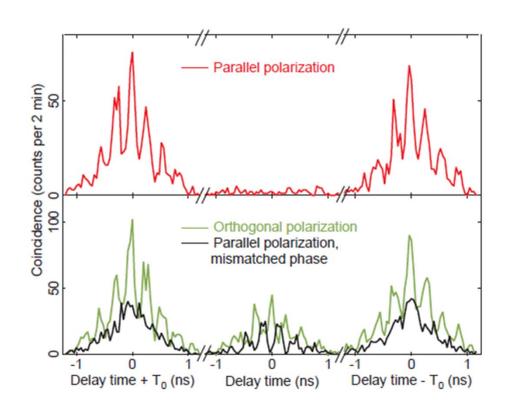
# Generation of a single-photon color-qubit from a QD exciton

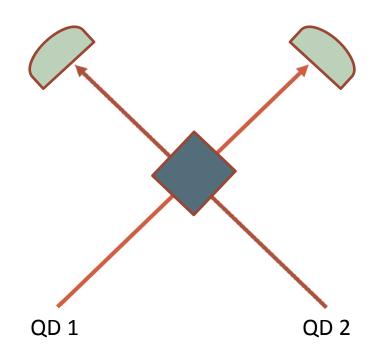
In a neutral QD, the elementary optical excitations are excitons (X0); the two linearly polarized exciton X0 lines are split due to electron-hole exchange by ~ 5 GHz



By controlling the pulse-shape, detuning and polarization of the resonant laser, we could generate a single-color photon or a two-color photonic qubit

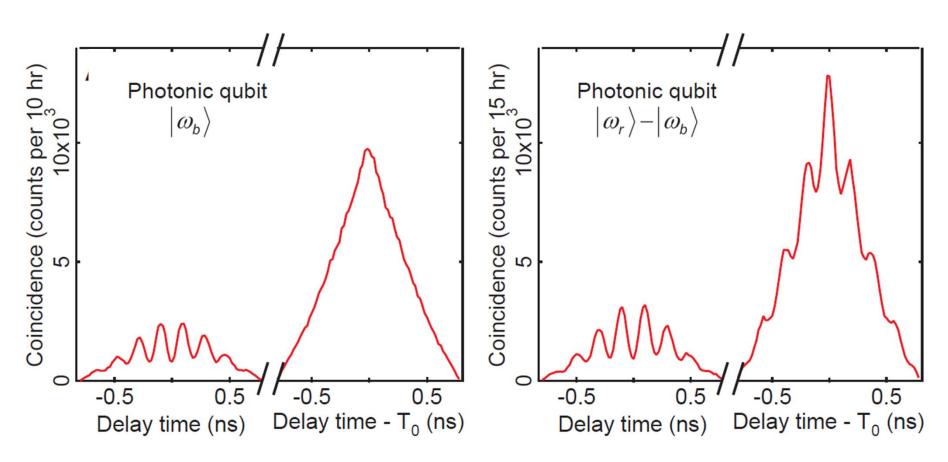
## Interference of photonic qubits (superposition of blue & red photons) coming from two neutral quantum dots





- For identical incident photons no coincidence detection (i.e. no counts around t = 0)
- 80% visibility in intereference of two photonic (color) qubits

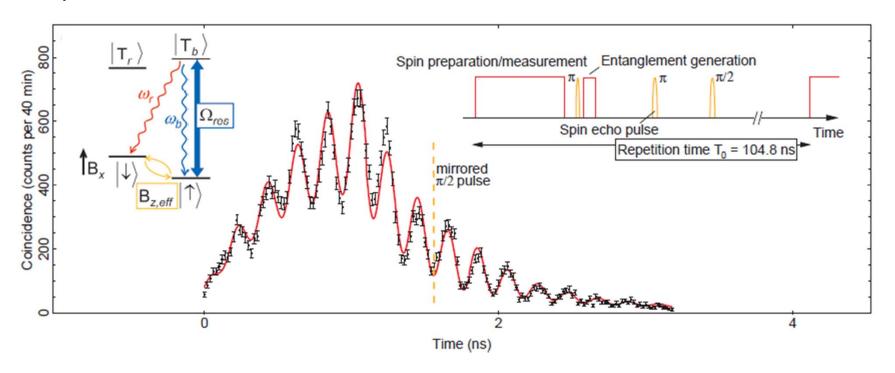
# Interference of one photon coming from a neutral QD with a photon from a single-electron-charged QD



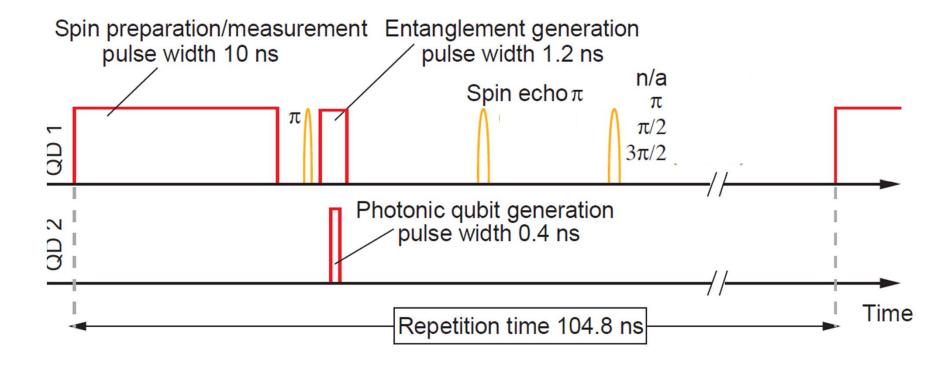
- Coincidences from the same-period are ideally ¼ of the other peaks, when the spin state is discarded.
- It is these coincidences that herald teleportation.

## Spin-photon entanglement with spin-echo

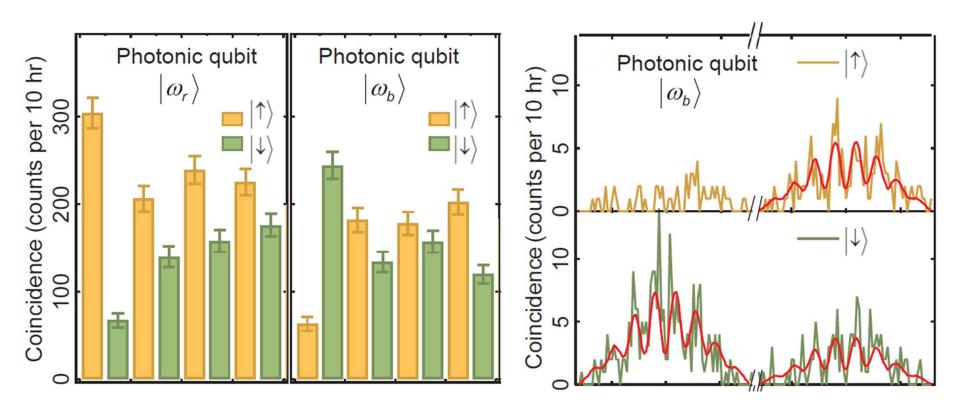
- To ensure that spin is detected after a coincidence event heralds successful teleportation, we need to prolong the spin-coherence
- Quantum correlations between the QD spin and the emitted singlephoton after an echo time of 13 ns



### Quantum teleportation pulse sequence

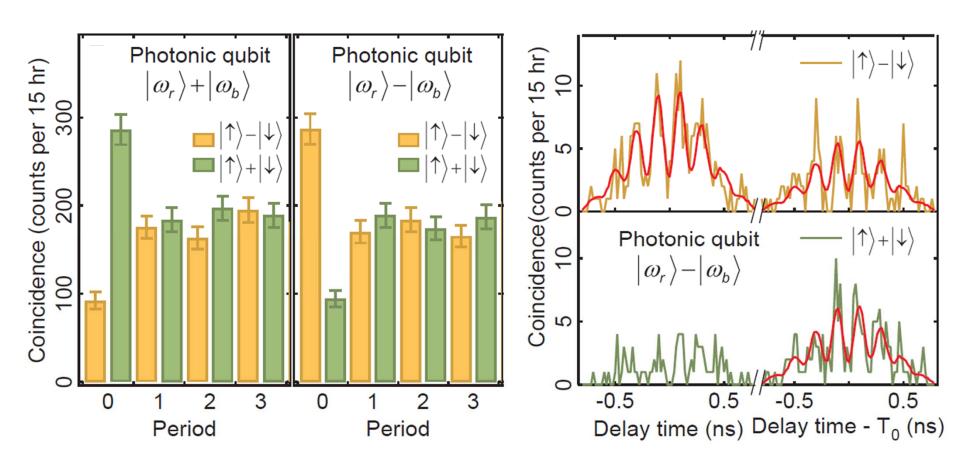


# «Classical teleportation» (3-fold coincidences)



• For a photonic qubit in  $|\omega_b\rangle$ , detection of a coincidence ensures that the photon from the entangled pair was in  $|\omega_r\rangle$  - which in turn fixes the spin to be in  $|\uparrow\rangle$ .

# Quantum correlations in teleportation (3-fold coincidences)



Overall teleportation fidelity:  $0.78 \pm 0.03$ 

### Outlook

- Spin-Spin entanglement
- Understanding and suppressing the role of hyperfine interactions.

#### Thanks to

- Weibo Gao
- Aymeric Delteil, Emre Togan, Parisa Fallahi, Javier Sanchez