

Recent Results from Single and Double Layer Quantum Hall Systems

Jim Eisenstein

Vaclav Cvicek

Aaron Finck

Debaleena Nandi

Jing Xia

Loren Pfeiffer

Ken West



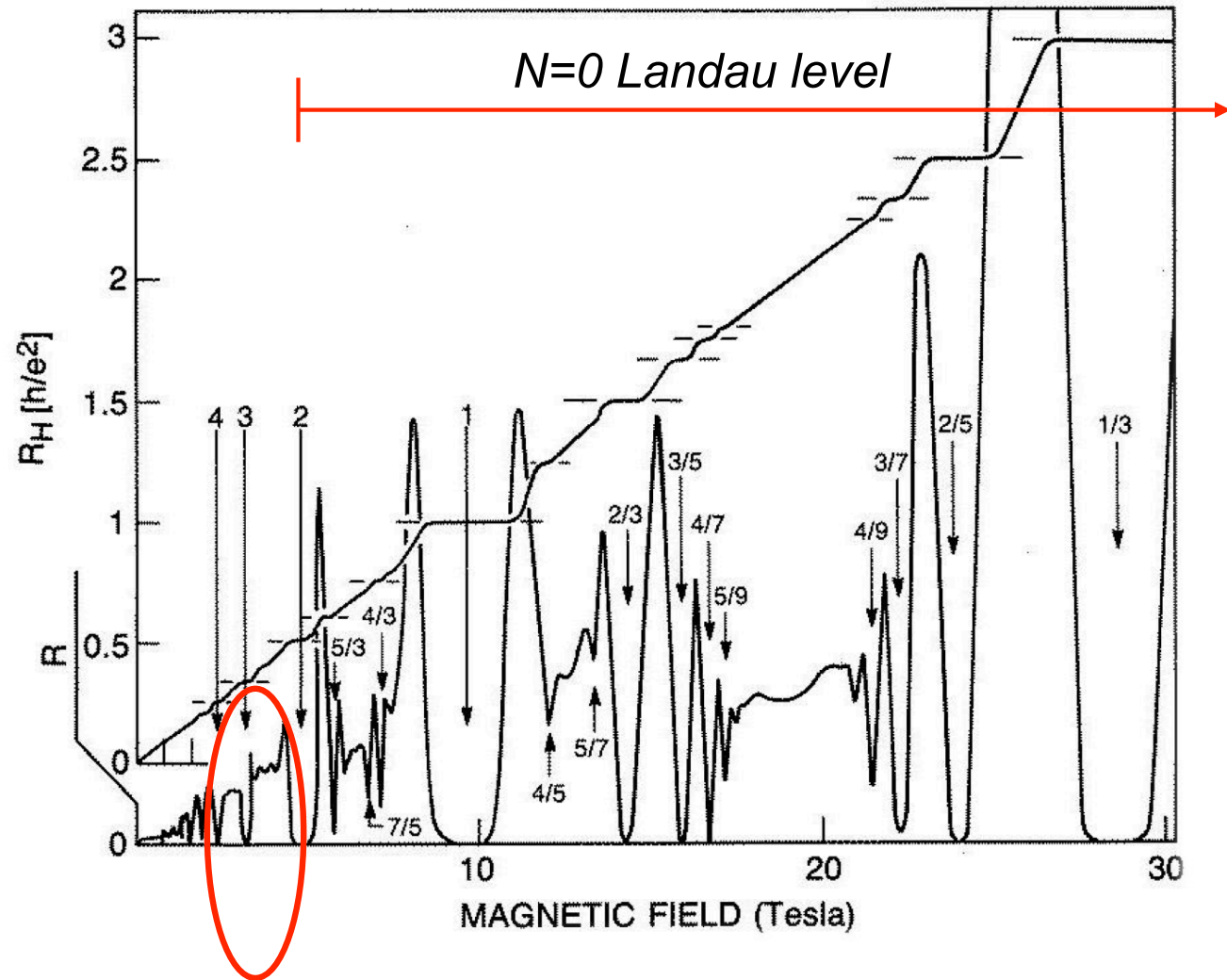
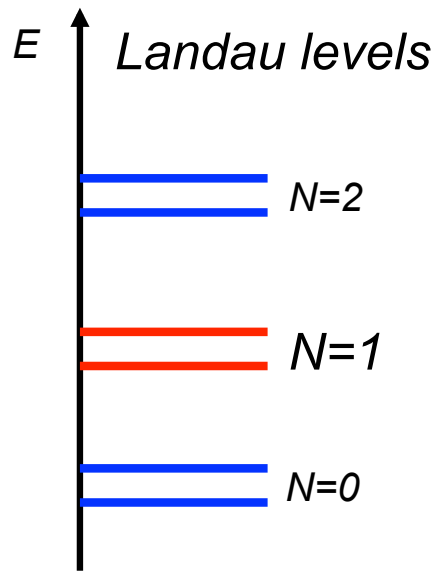
U.S. DEPARTMENT OF ENERGY: BASIC ENERGY SCIENCE



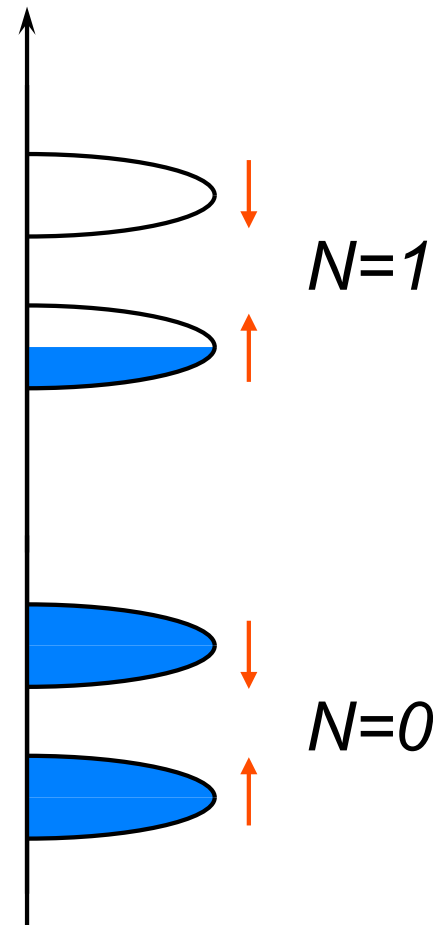
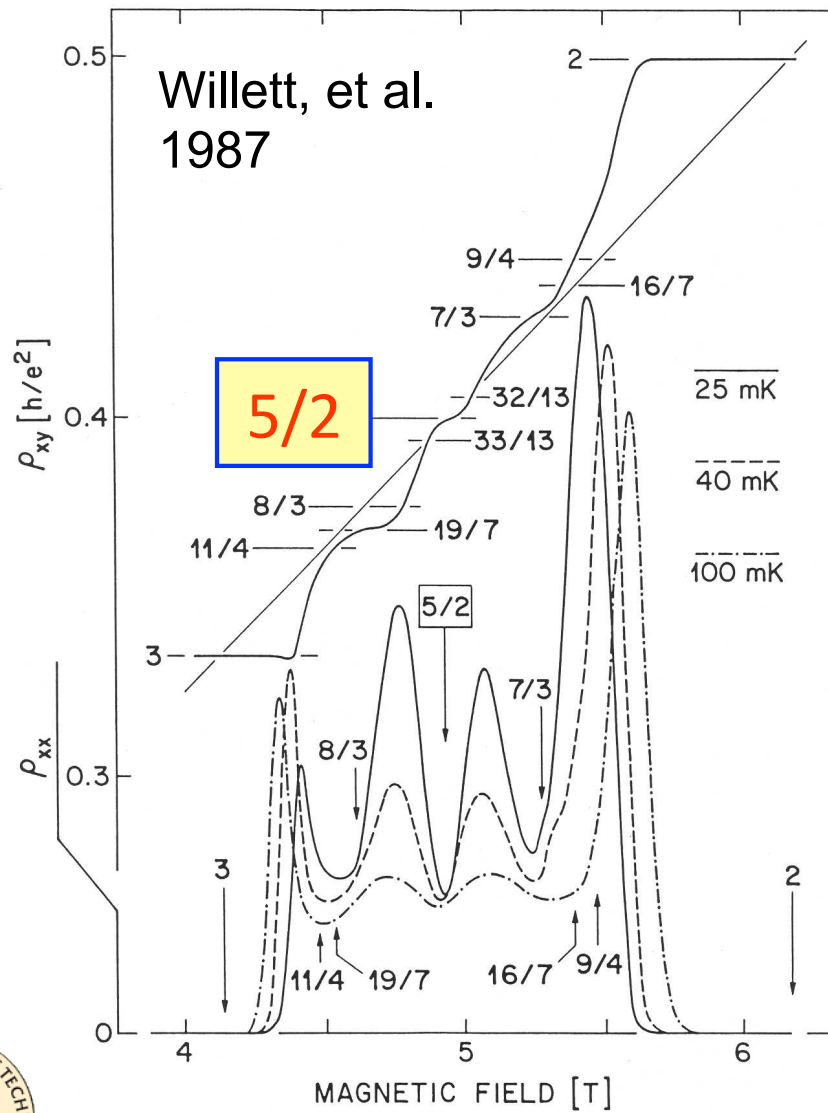
Division of Materials Sciences

Single layer FQHE systems

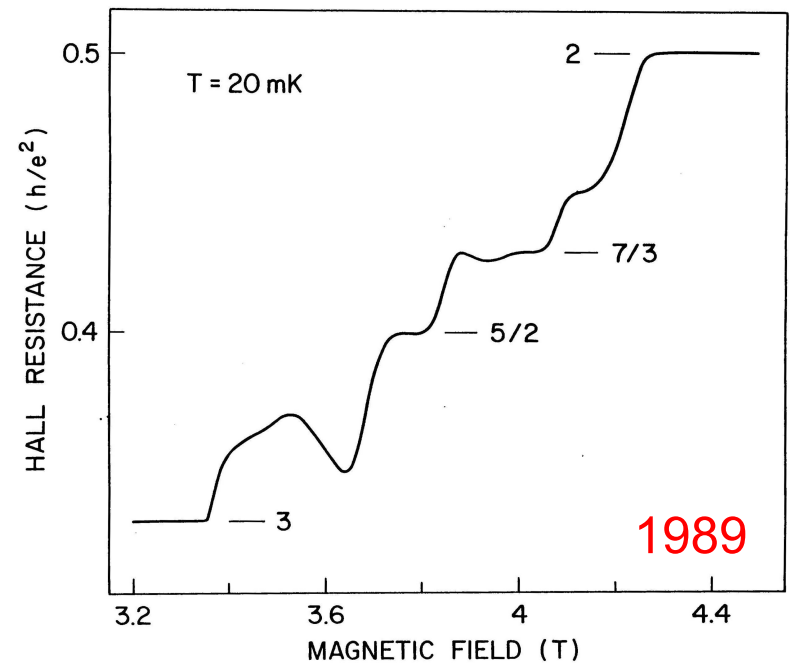
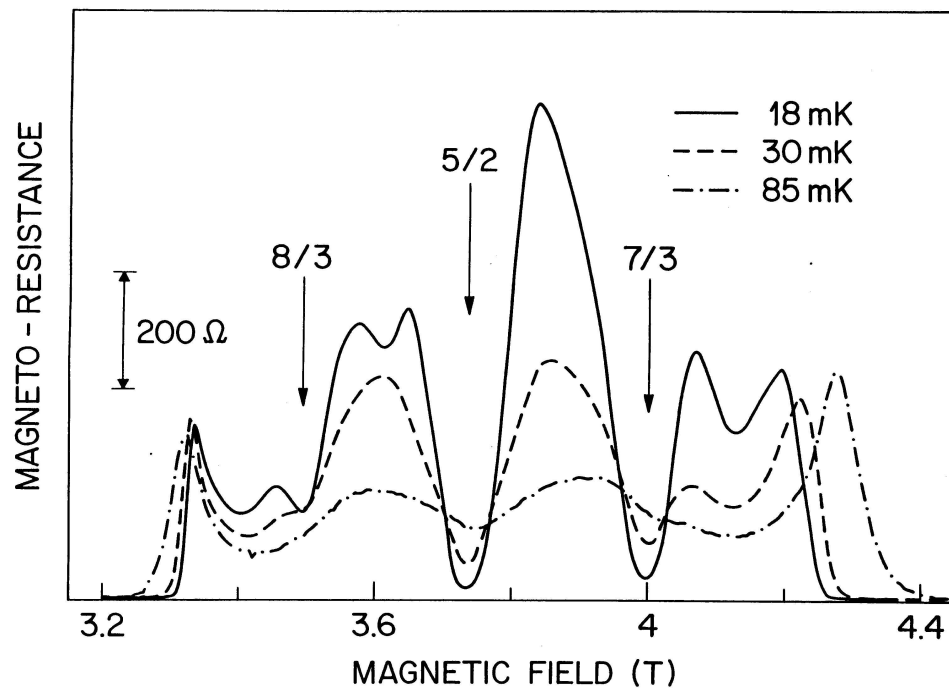
$N=1$ Landau level



Even-denominator FQHE



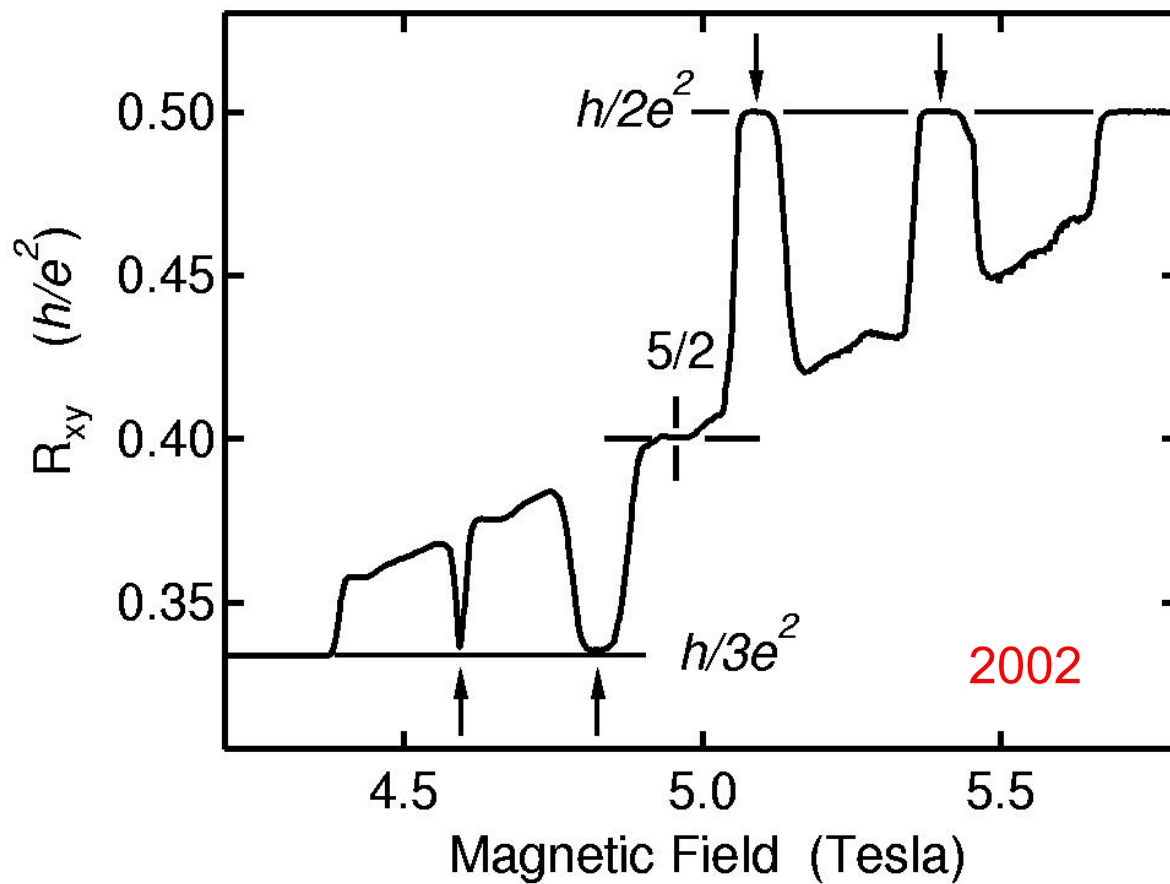
Better sample



Vanishing R_{xx} and good Hall plateau



Re-entrant integer QHE states



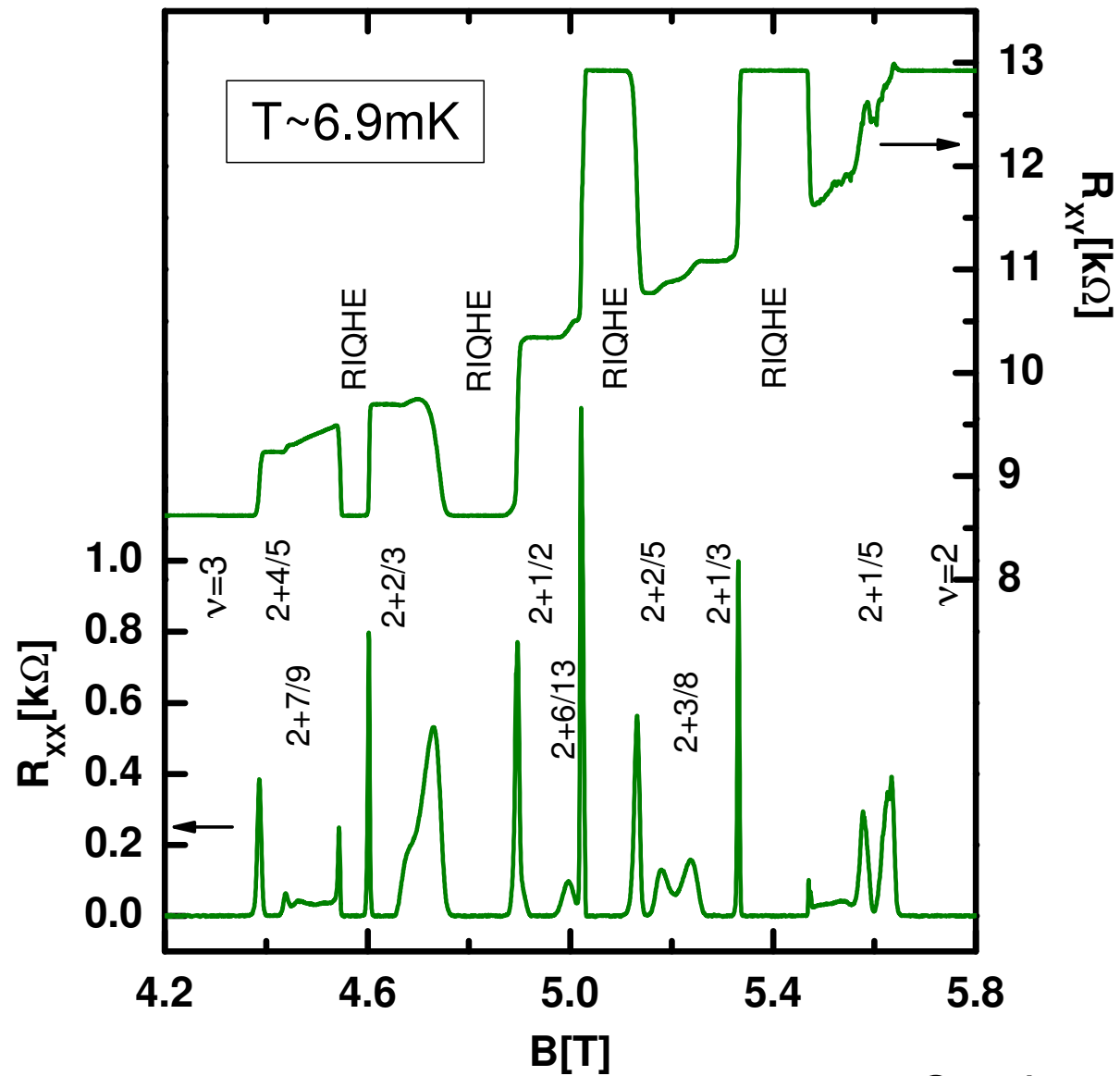
Integer Hall quantization near
 $\nu - 2 \approx 2/7, 3/7, 4/7, \text{ and } 5/7$



New Collective Insulators

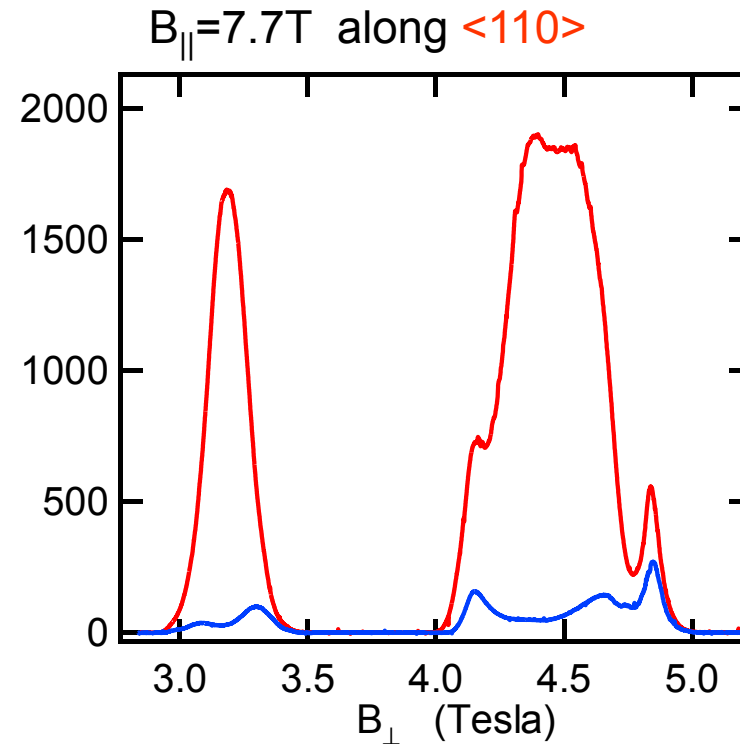
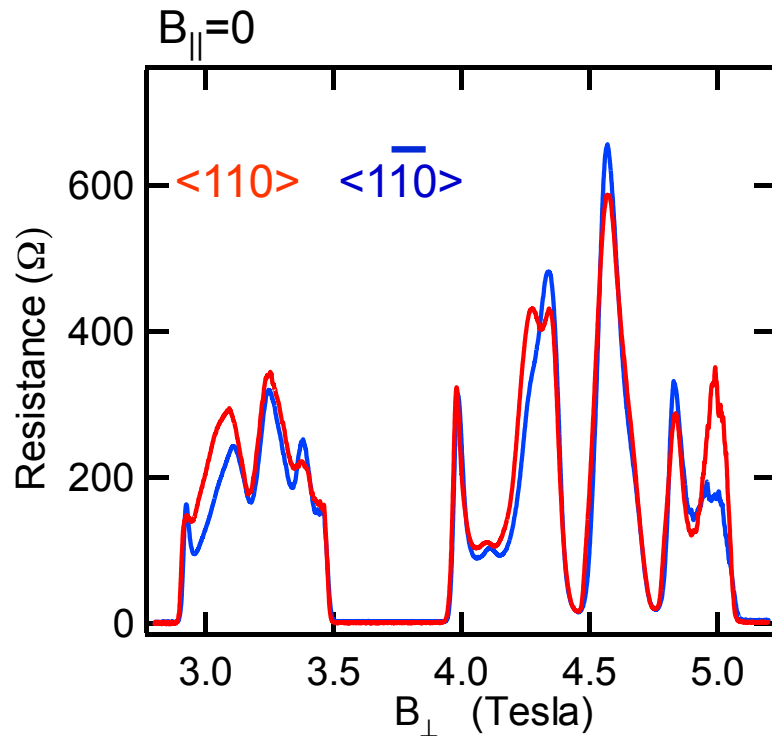


Best data yet?



Csathy, et al PRL 2010

5/2 State in tilted fields



*Tilting destroys 5/2 and 7/2 FQHE states **and** produces strongly anisotropic transport. Competing phases must be nearby.*



Nearly degenerate phases at 5/2

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Incompressible Paired Hall State, Stripe Order, and the Composite Fermion Liquid Phase in Half-Filled Landau Levels

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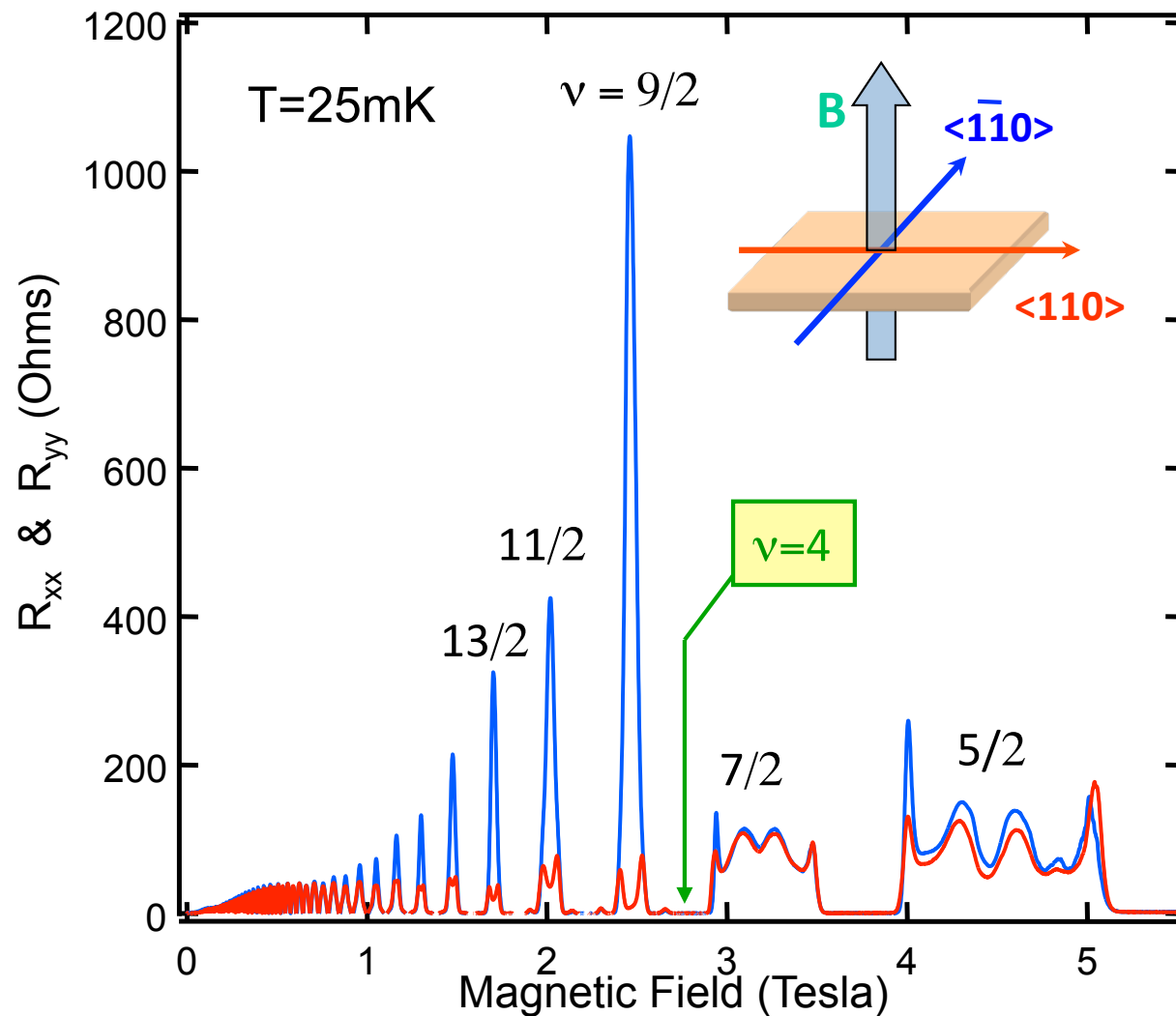
(Received 14 June 1999)

We consider the two lowest Landau levels at half filling. In the higher Landau level ($\nu = 5/2$), we find a first-order phase transition separating a compressible striped phase from a paired quantum Hall state, which is identified as the Moore-Read state. The critical point is very near the Coulomb potential and the transition can be driven by increasing the width of the electron layer. We find a much weaker transition (either second-order or a crossover) from pairing to the composite fermion Fermi-liquid behavior. A very similar picture is obtained for the lowest Landau level, but the transition point is not near the Coulomb potential.

*Wider wells stabilize FQHE at 5/2.
Narrowing well drives transition to stripe phase.
Tilting narrows wavefunction?*

cf. Peterson, Jolicoeur, and Das Sarma, 2008

Transport anisotropy in $N \geq 2$ Landau levels



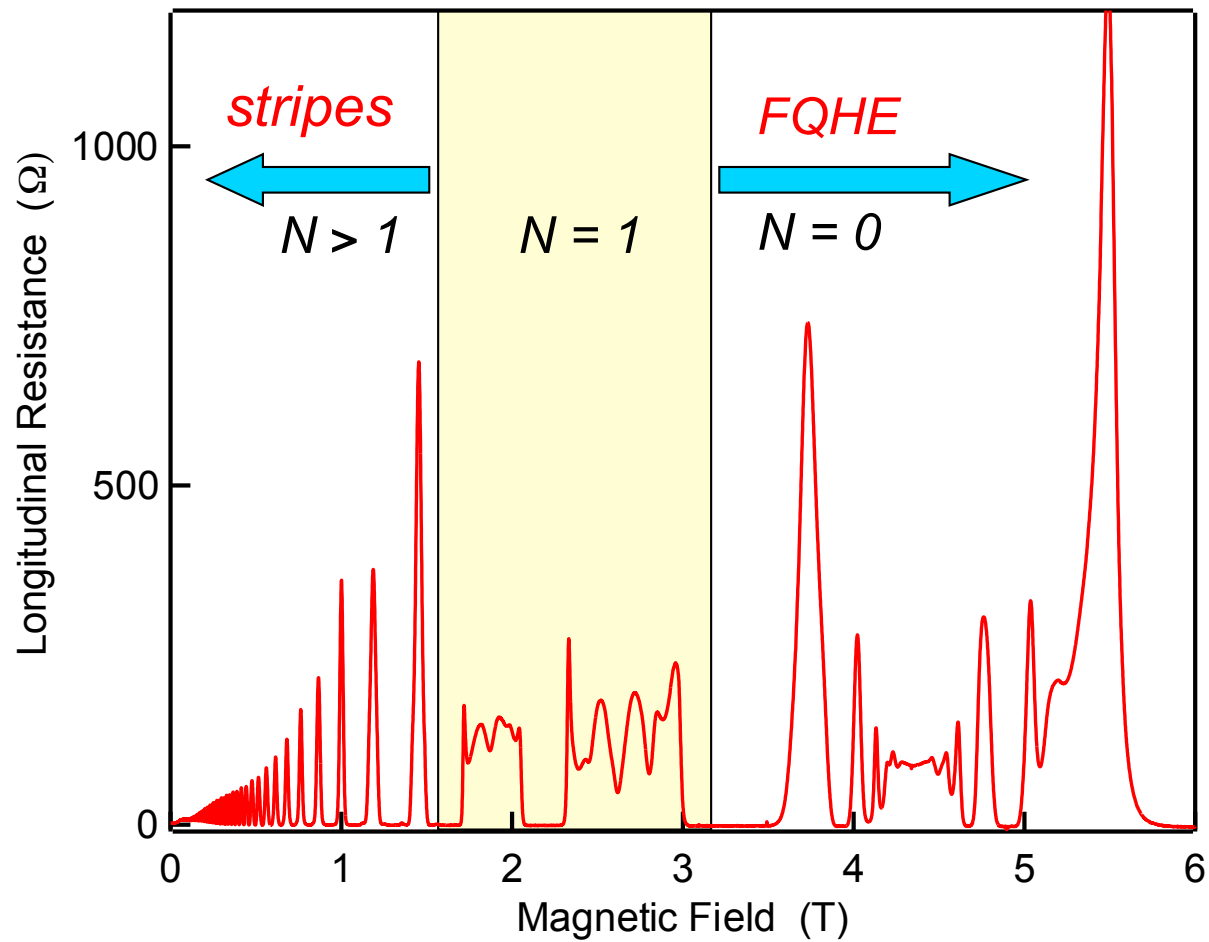
Evidence for electron stripe phases.



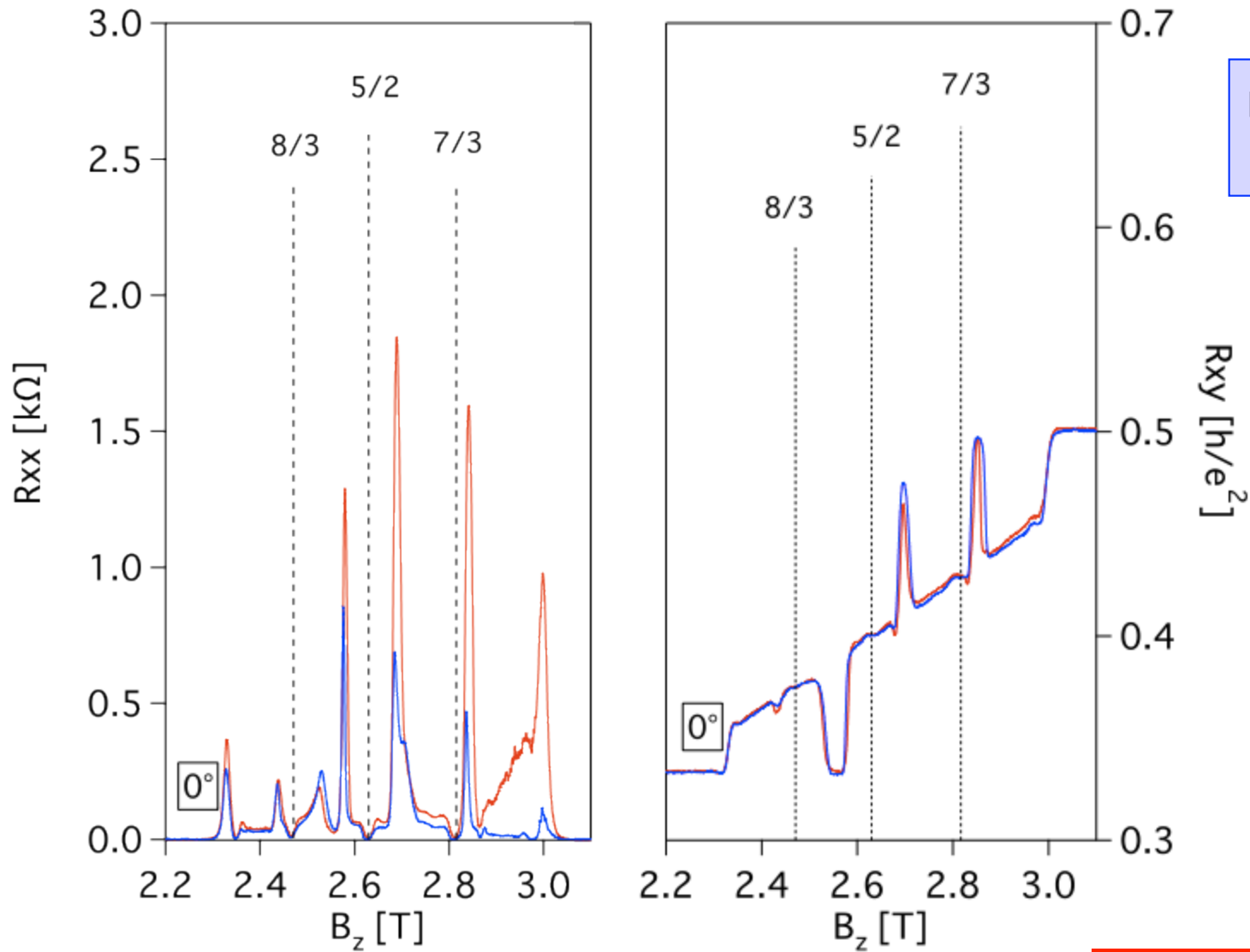
$N = 2, 3, \dots$

$N = 0 \text{ \& \; } 1$

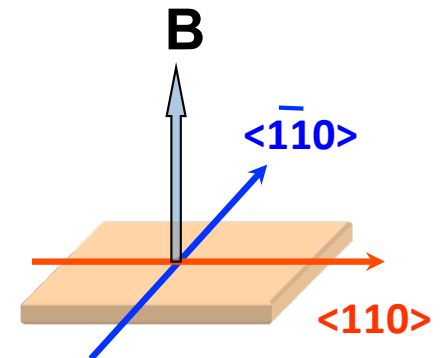
$N=1$ LL: Land of many phases!



400 Å quantum well

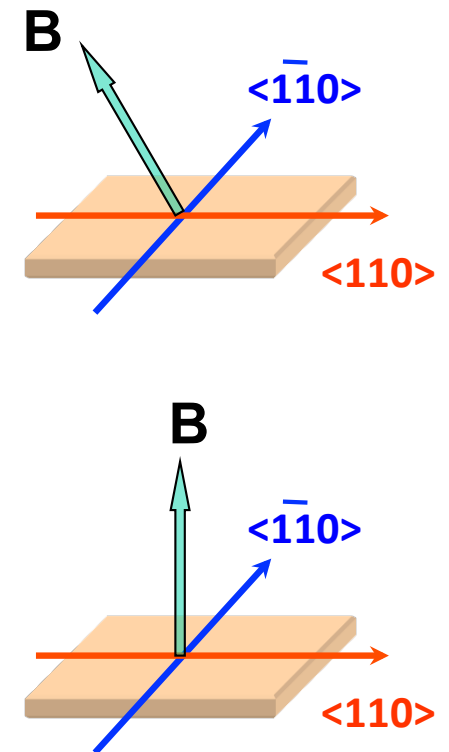
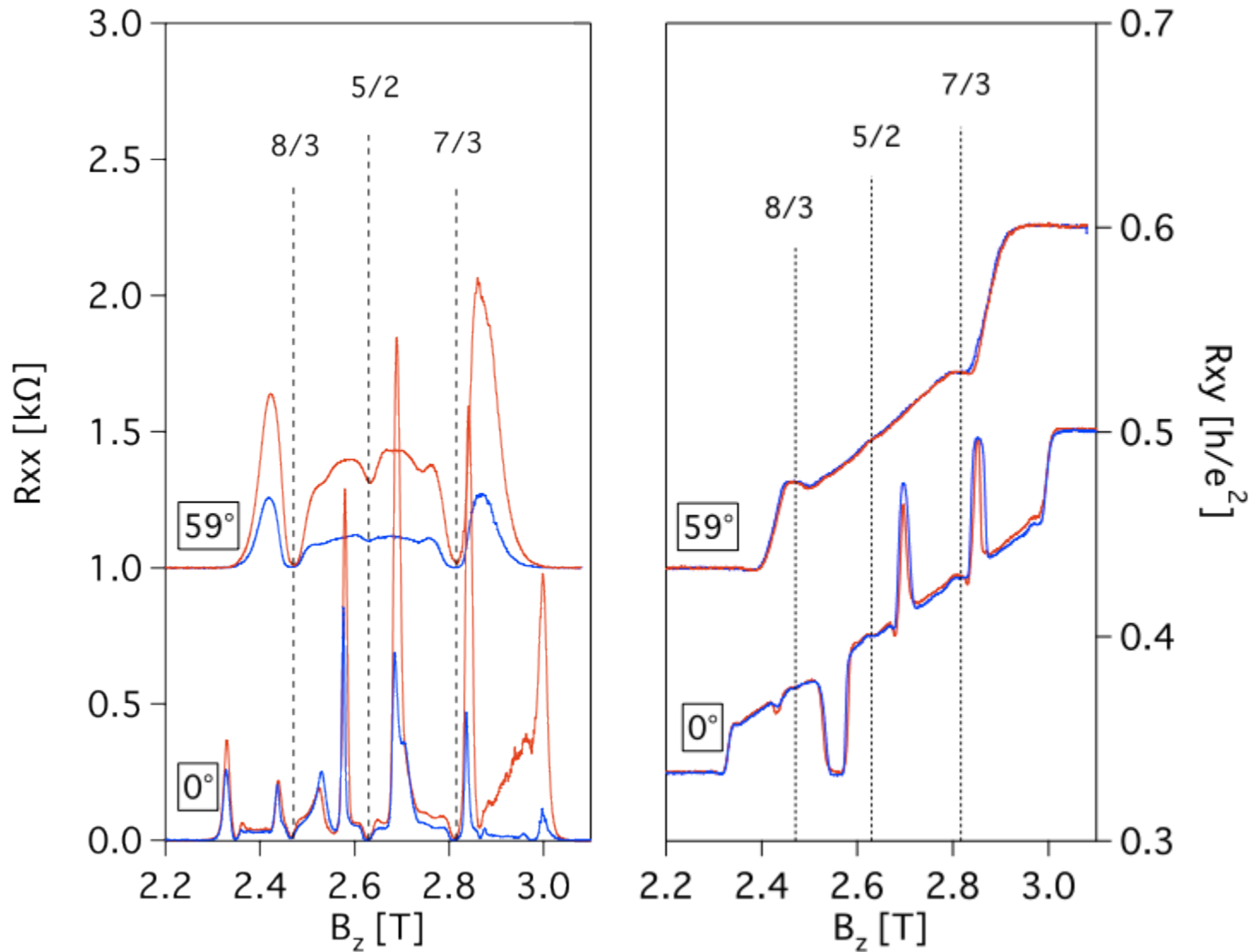


$$n = 1.6 \times 10^{11} \text{ cm}^{-2}$$
$$\mu = 16 \text{ M}$$



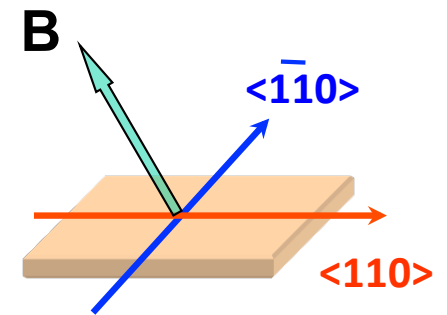
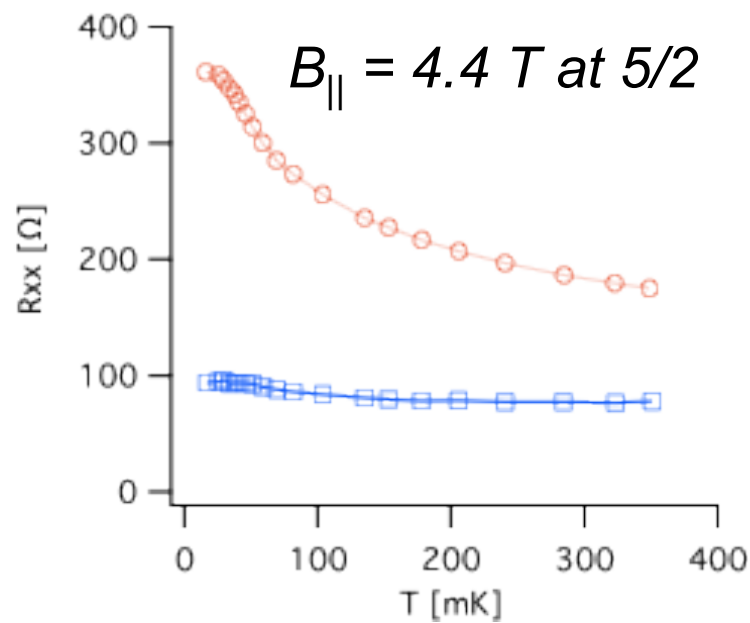
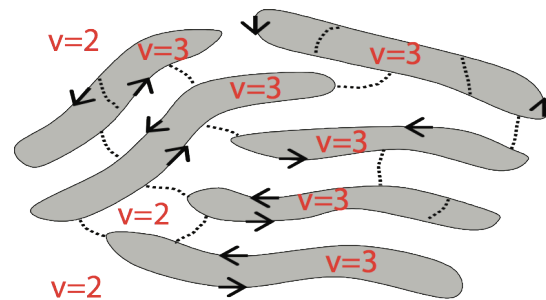
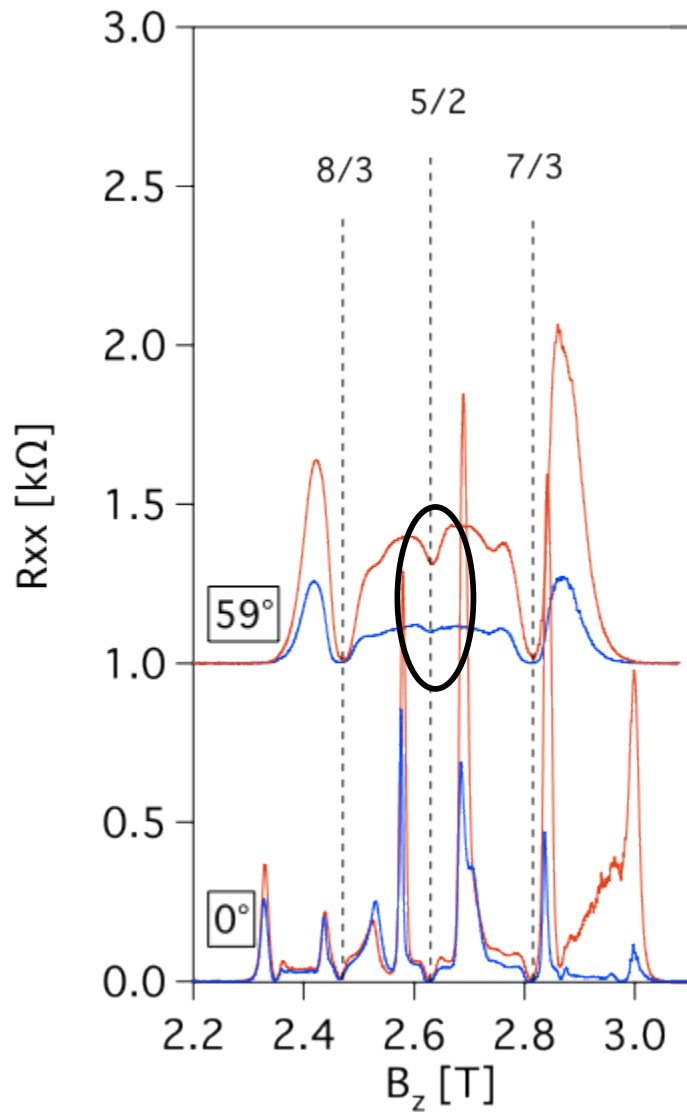
Familiar phenomenology

400 Å quantum well



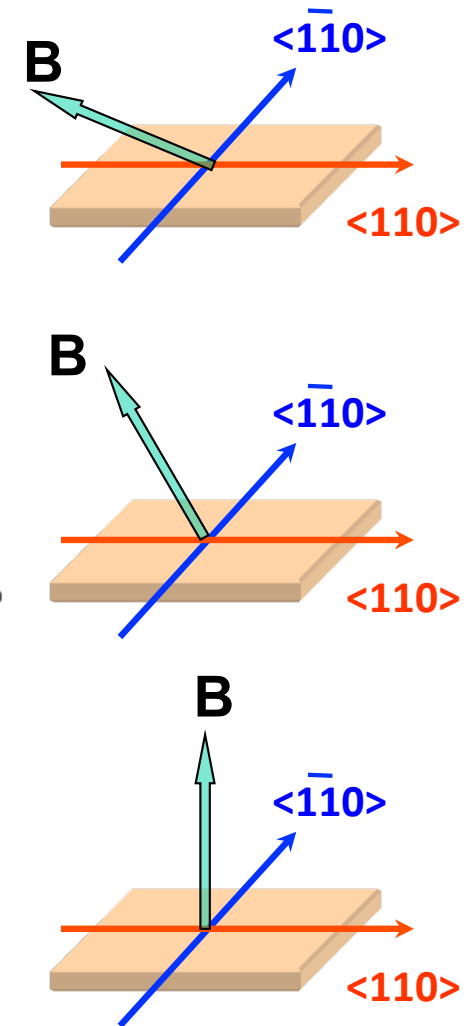
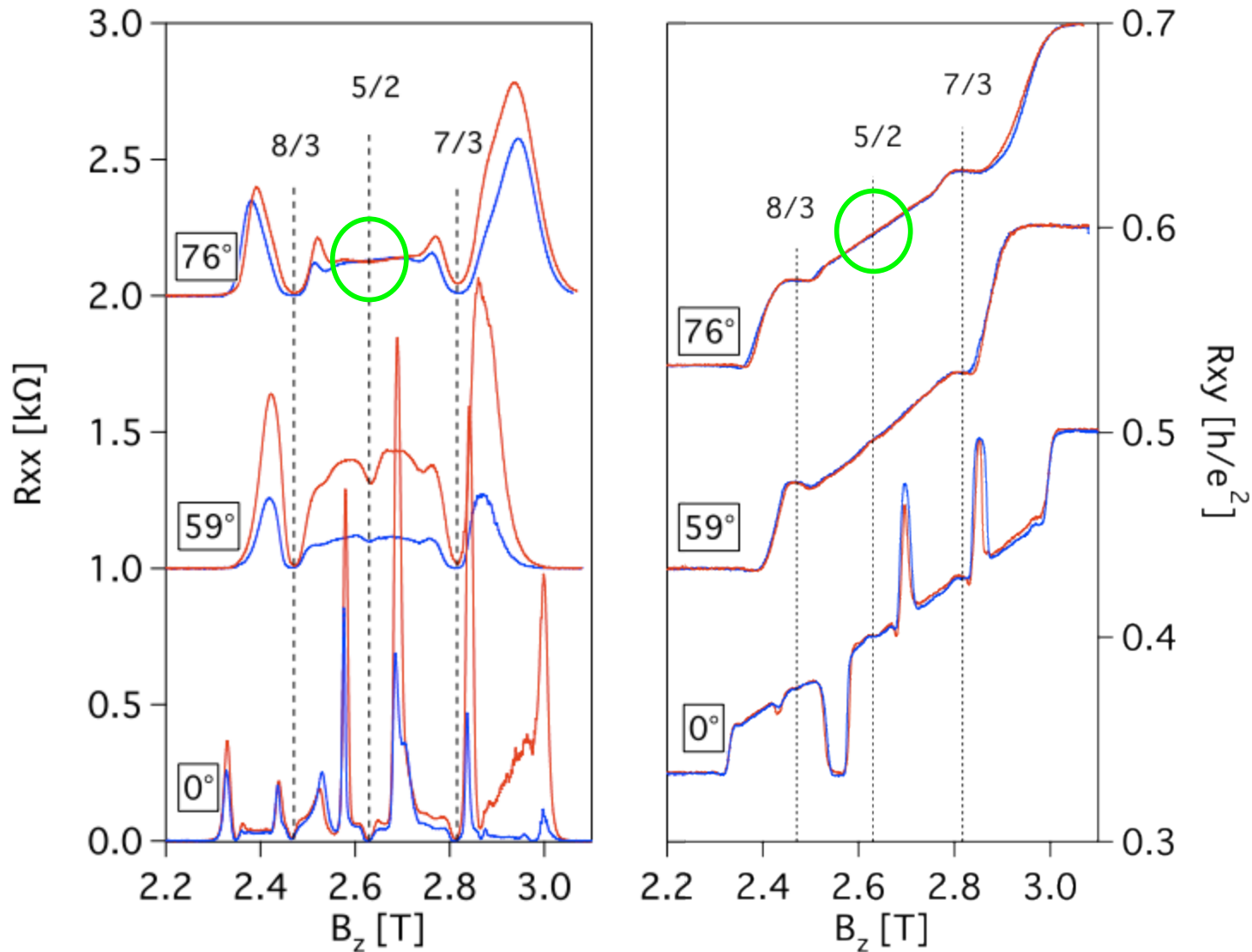
5/2 FQHE and RIQHE destroyed

400 Å quantum well



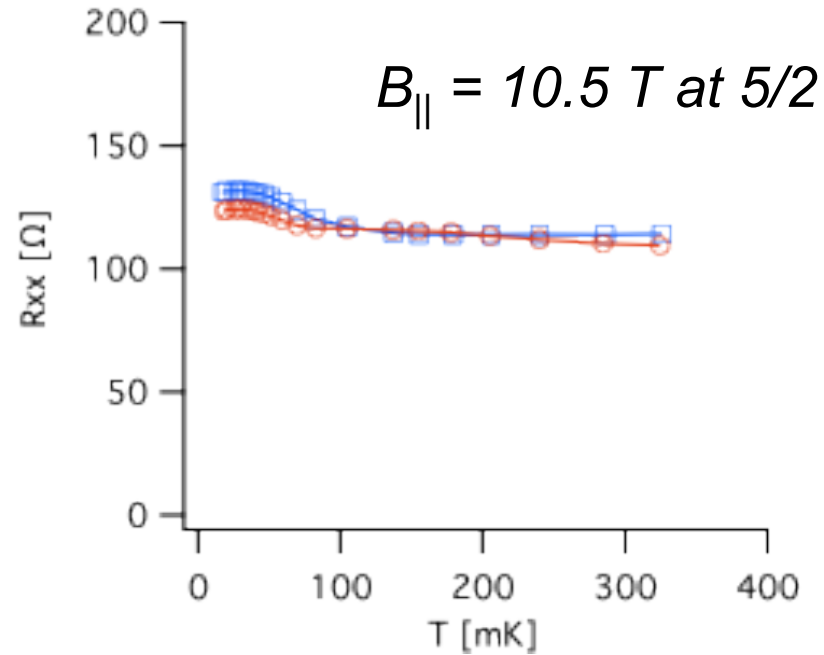
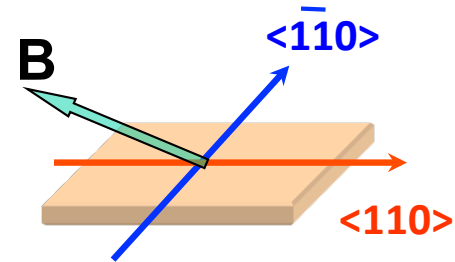
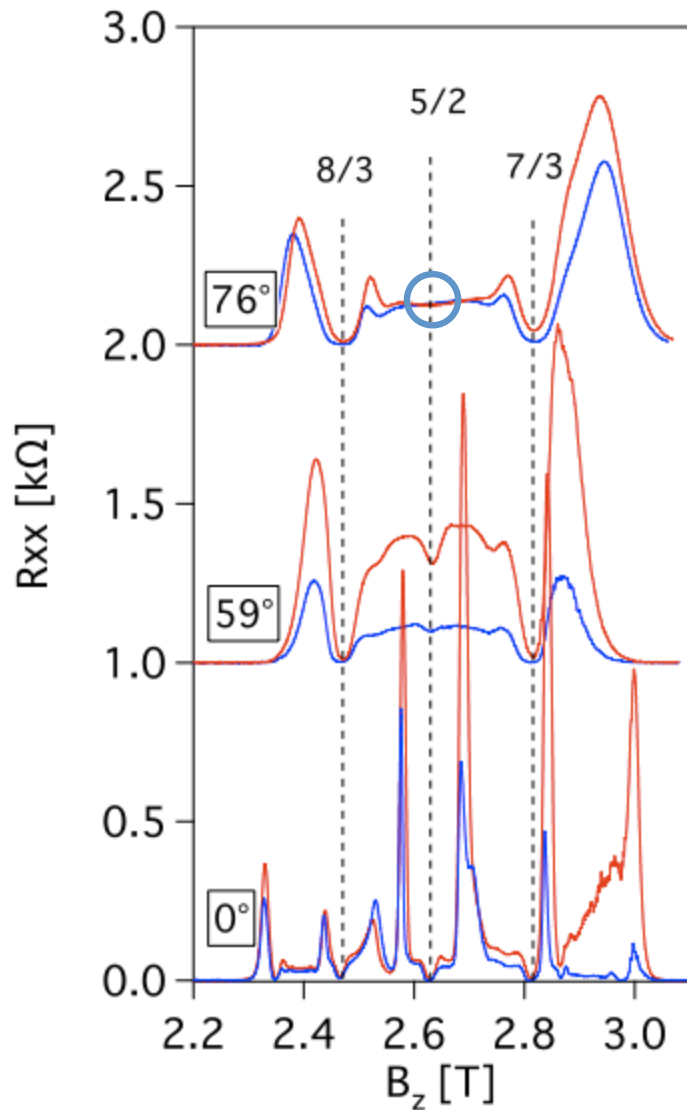
Stripe or Nematic at 5/2

400 Å quantum well



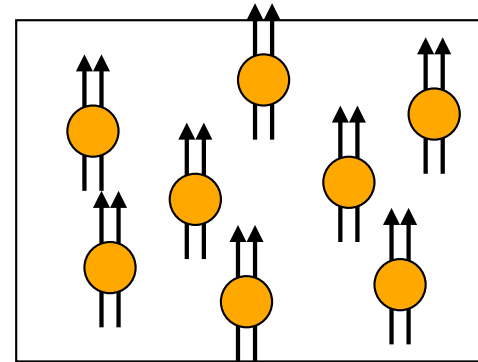
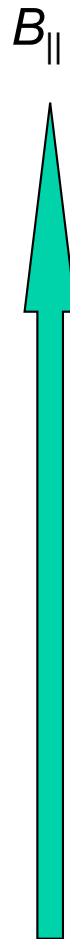
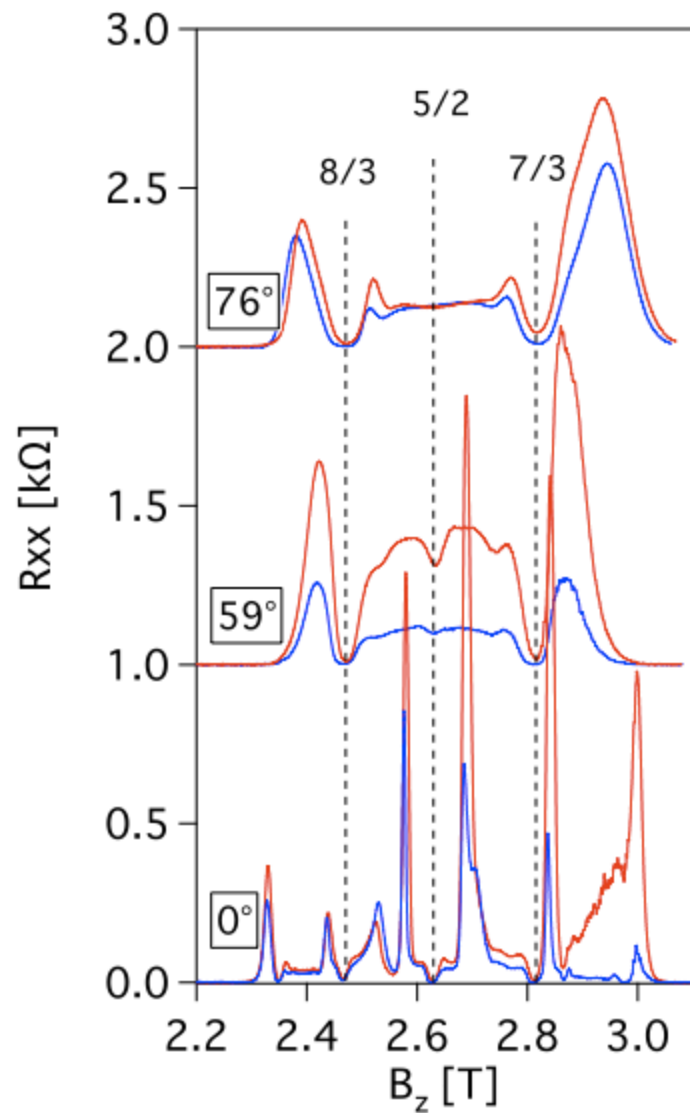
Re-entrant *Isotropic* Compressible (RIC) Phase at $5/2$!

400 Å quantum well

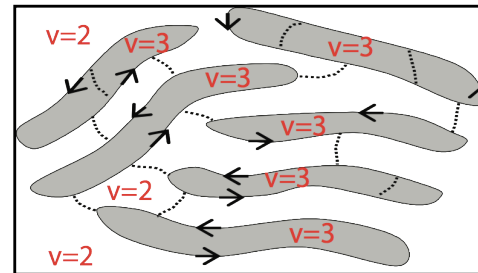


Composite Fermi Sea at $5/2$?

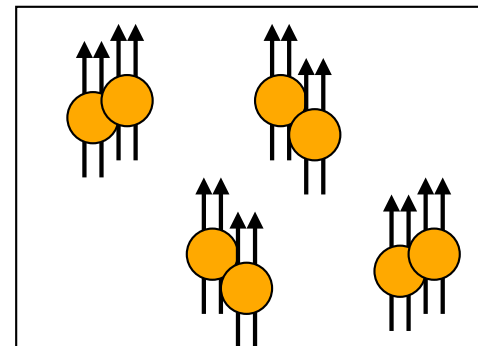
Quantum zoology at $\nu = 5/2$



*CF
Fermi Liquid*



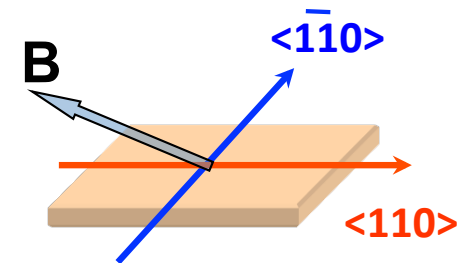
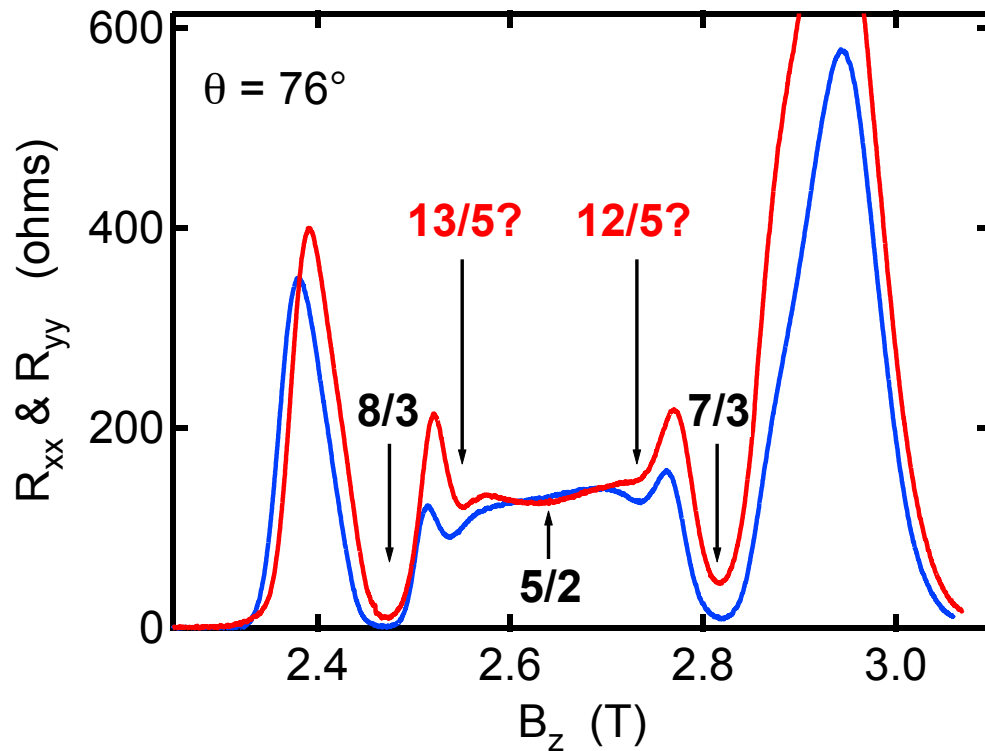
*Stripes or
Nematic*



*Paired CF
FQHE*

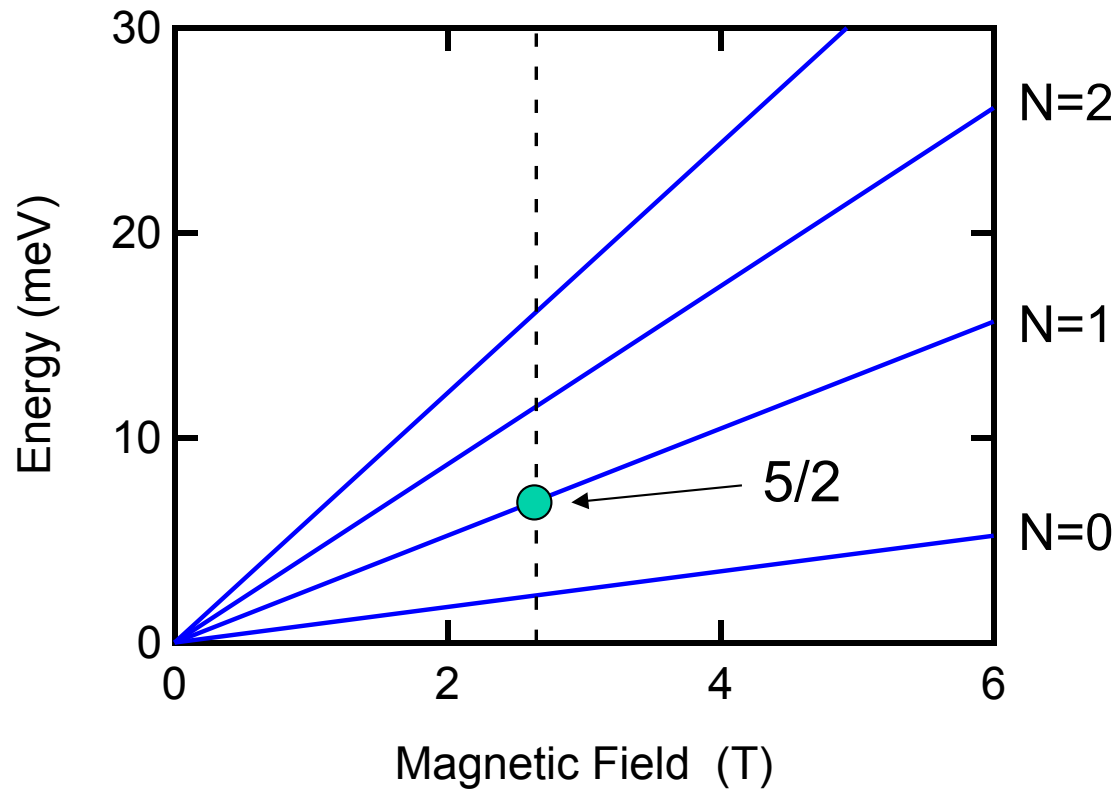
What drives these transitions?

400 Å Quantum Well

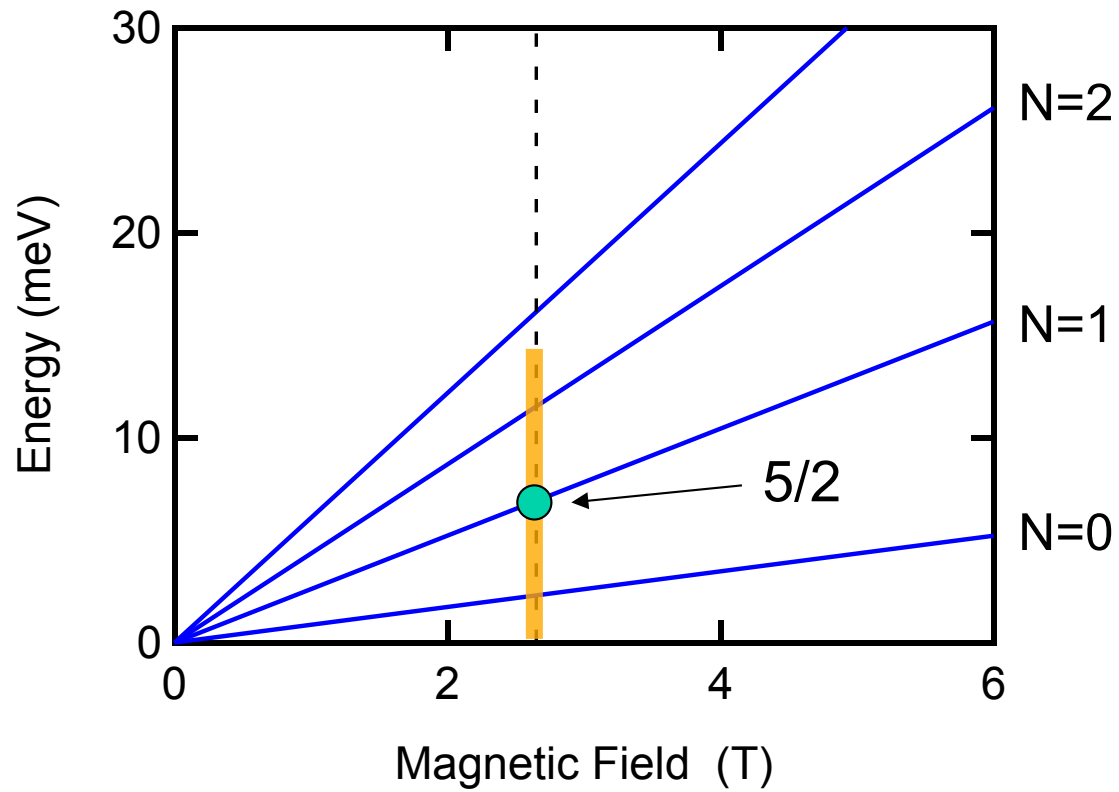


Looks like the lowest Landau level

Landau level mixing

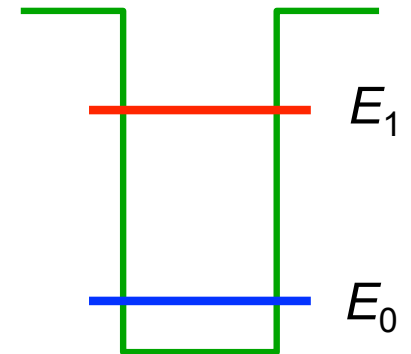
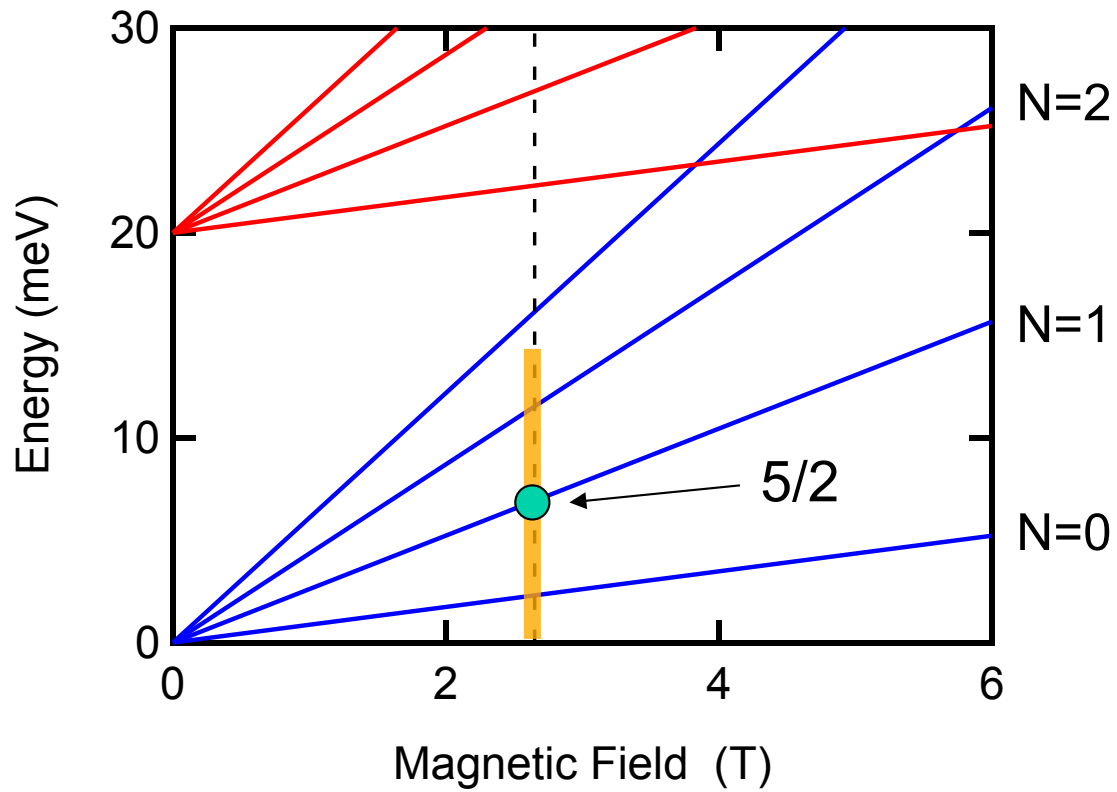


Landau level mixing



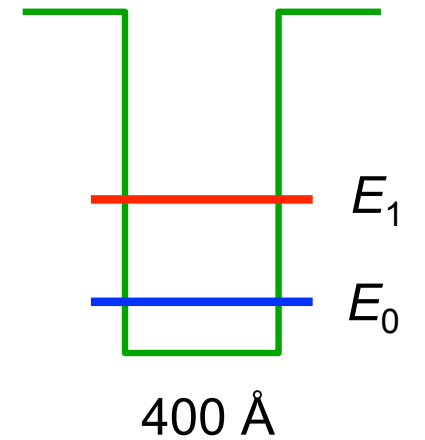
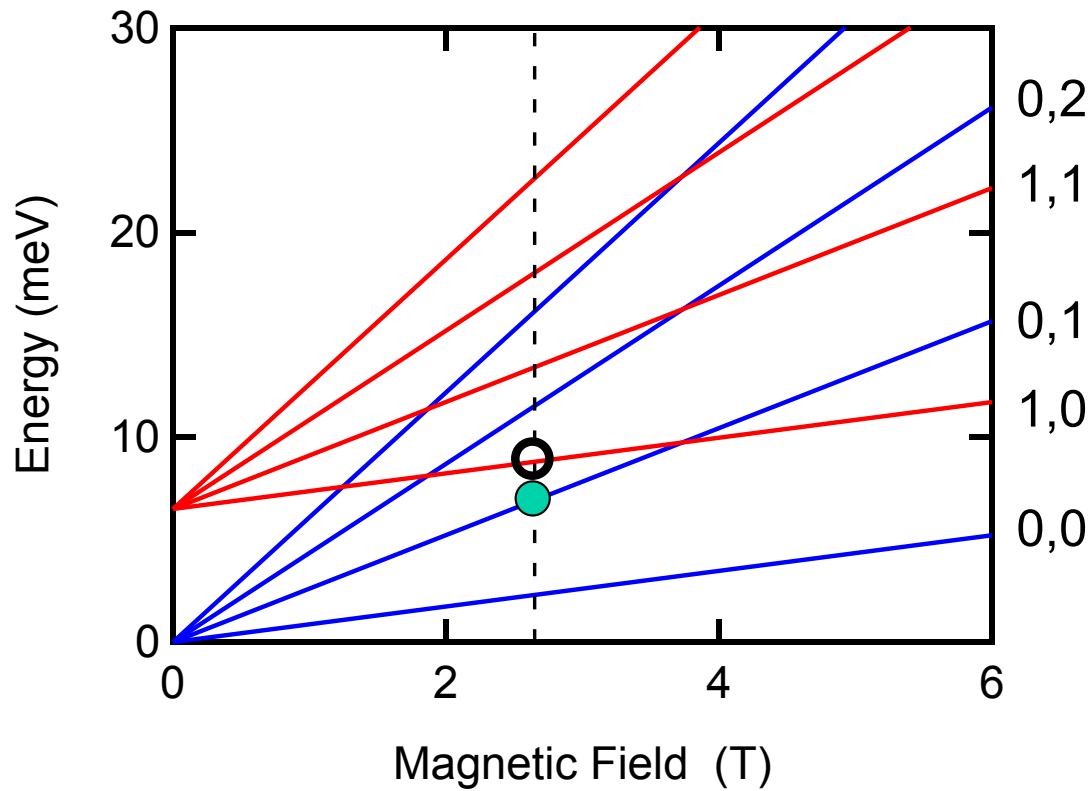
$$\frac{e^2}{\epsilon \ell} \approx 1.5 \hbar \omega_c \quad @ B = 2.65 \text{ T}$$

Landau level mixing



Second subband usually ignored.

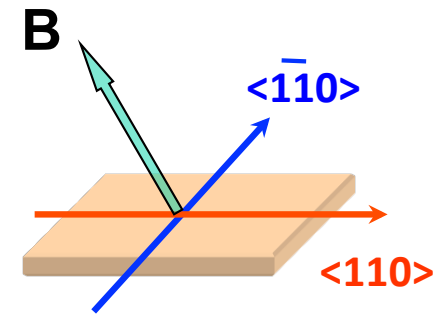
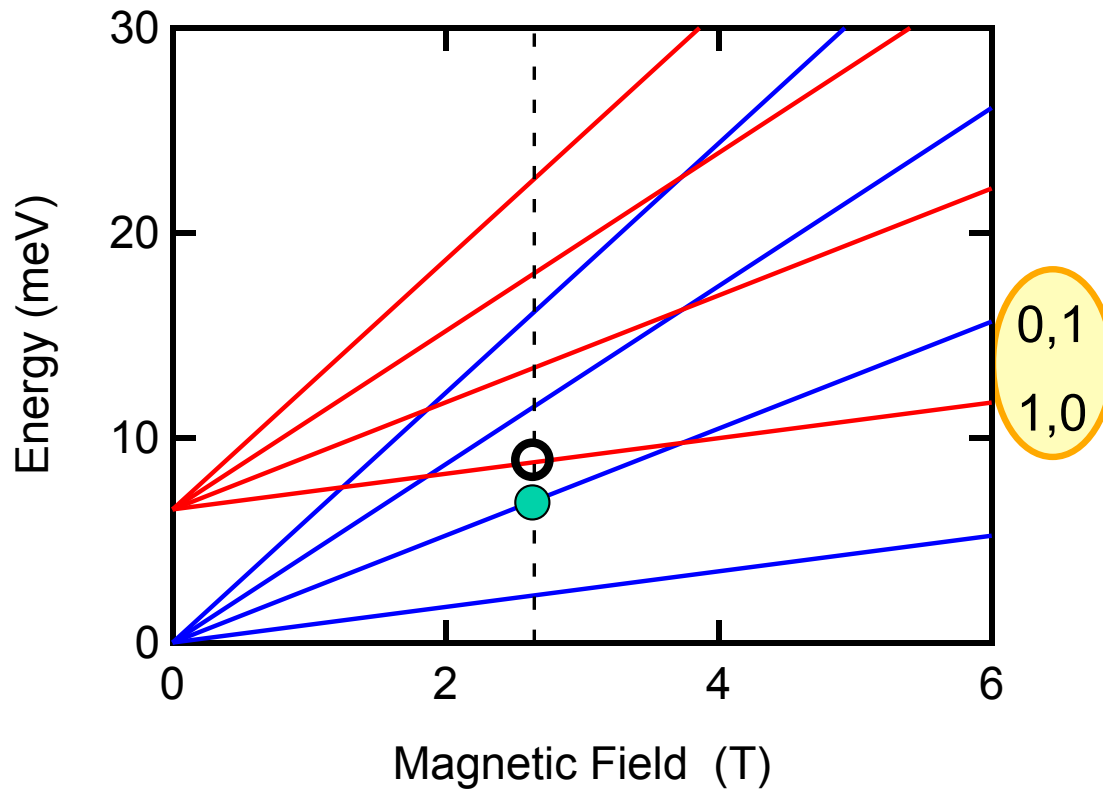
Subband – Landau level mixing



$$\Delta E_{10} = 6.5 \text{ meV}$$

Nearest empty level is $N=0$ LL in *second* subband

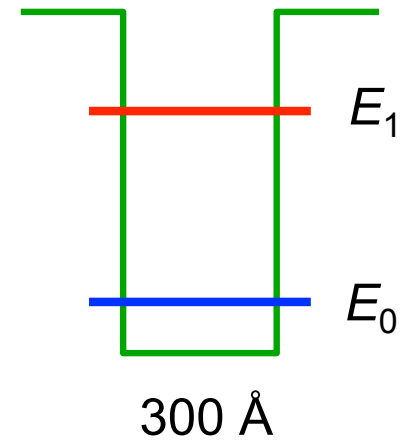
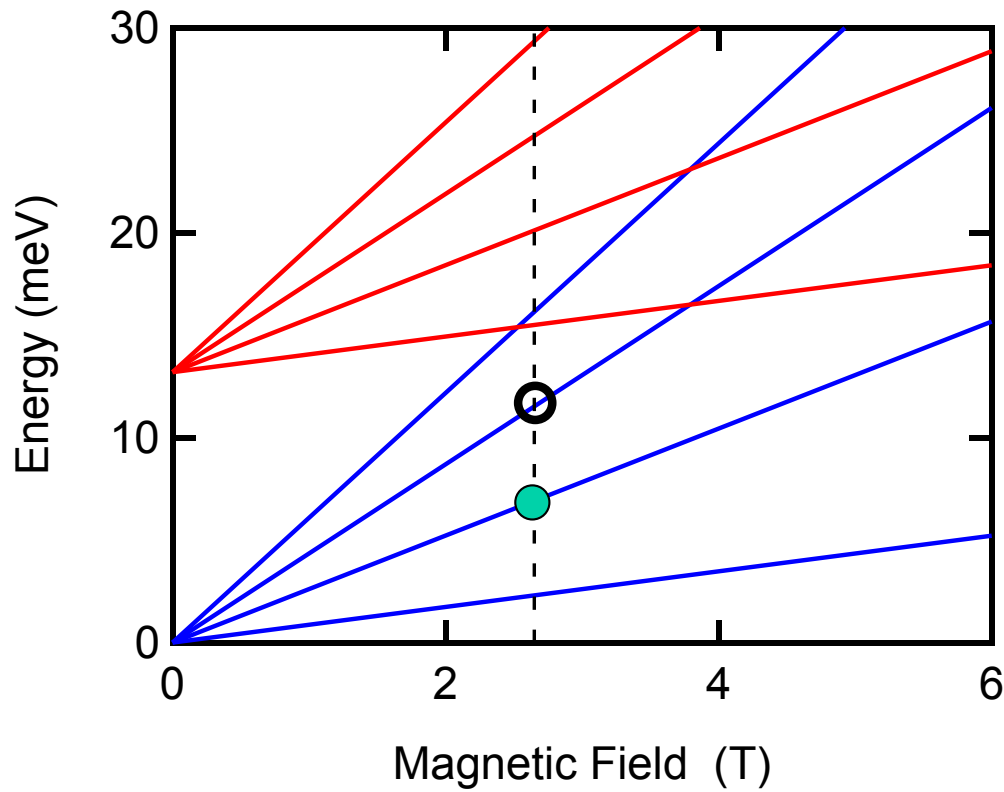
Subband – Landau level mixing



$$H' \sim B_{\parallel} z p_y$$

Tilting induces SB-LL coupling at the single particle level

Test: Narrower quantum well sample

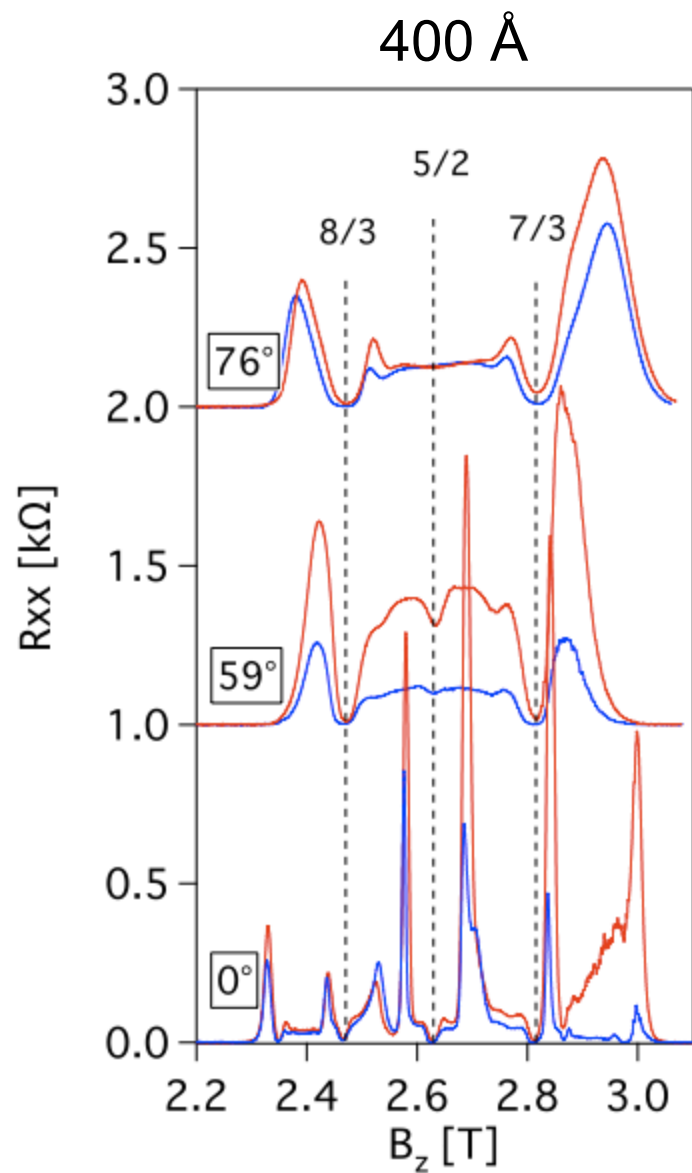


$$\Delta E_{10} = 13.2 \text{ meV}$$

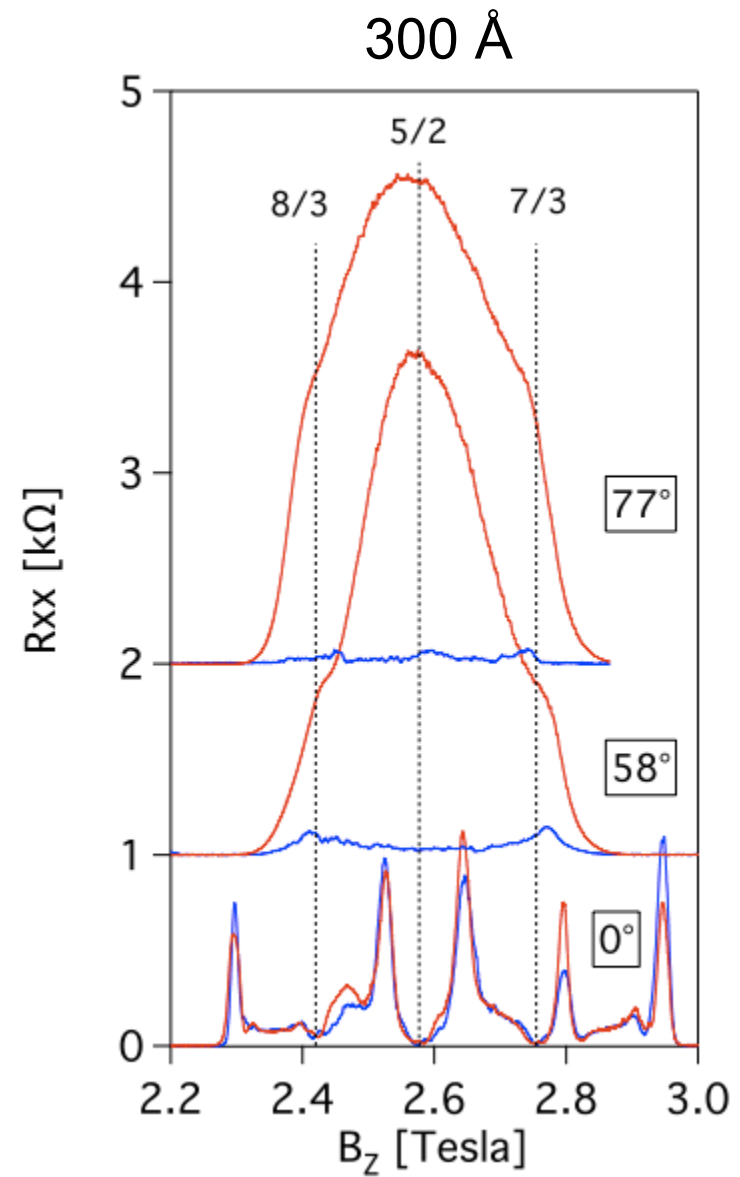
$$n = 1.6 \times 10^{11} \text{ cm}^{-2}$$
$$\mu = 16 \text{ M}$$

*Nearest empty level is $N=2$ LL in **ground** subband*

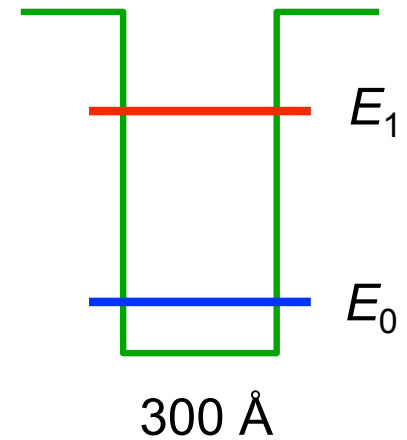
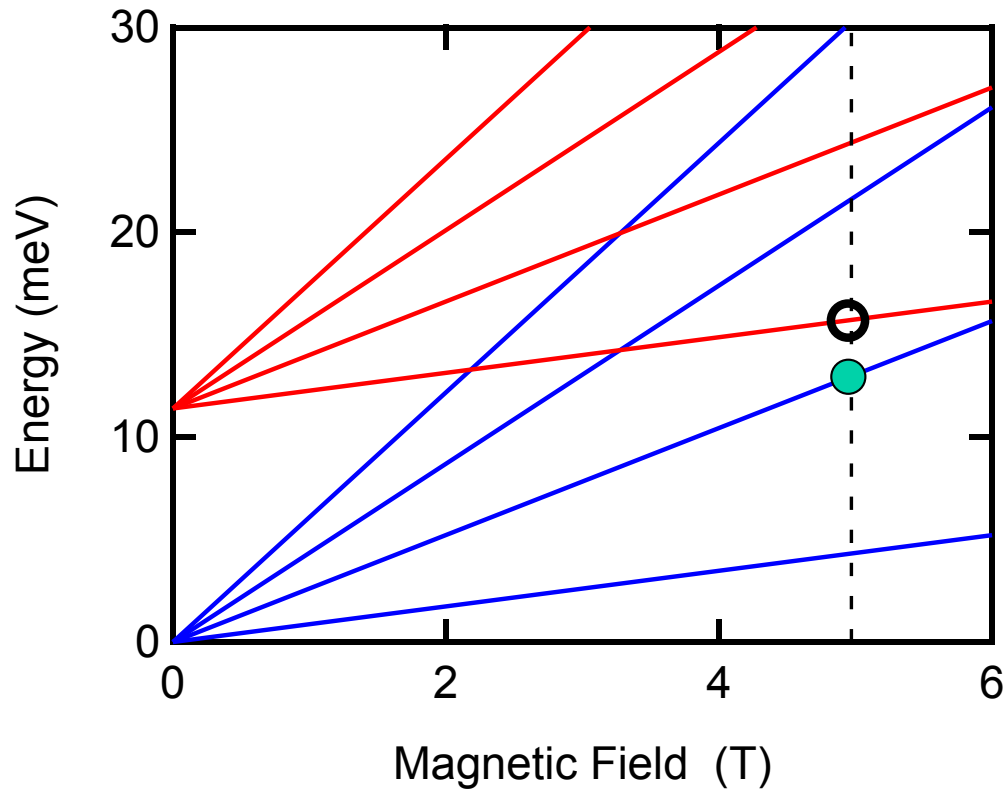
Dramatic difference between samples



$B_{||}$



Is well width the key difference?

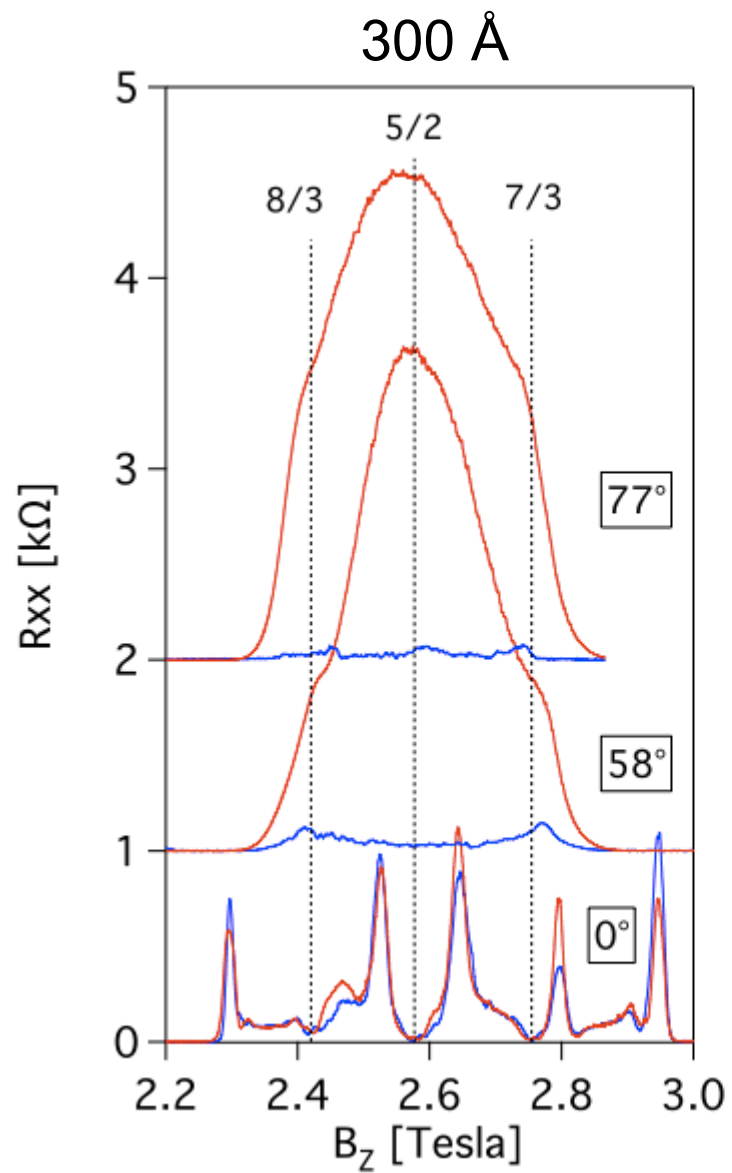


$$\Delta E_{10} = 11.4 \text{ meV}$$

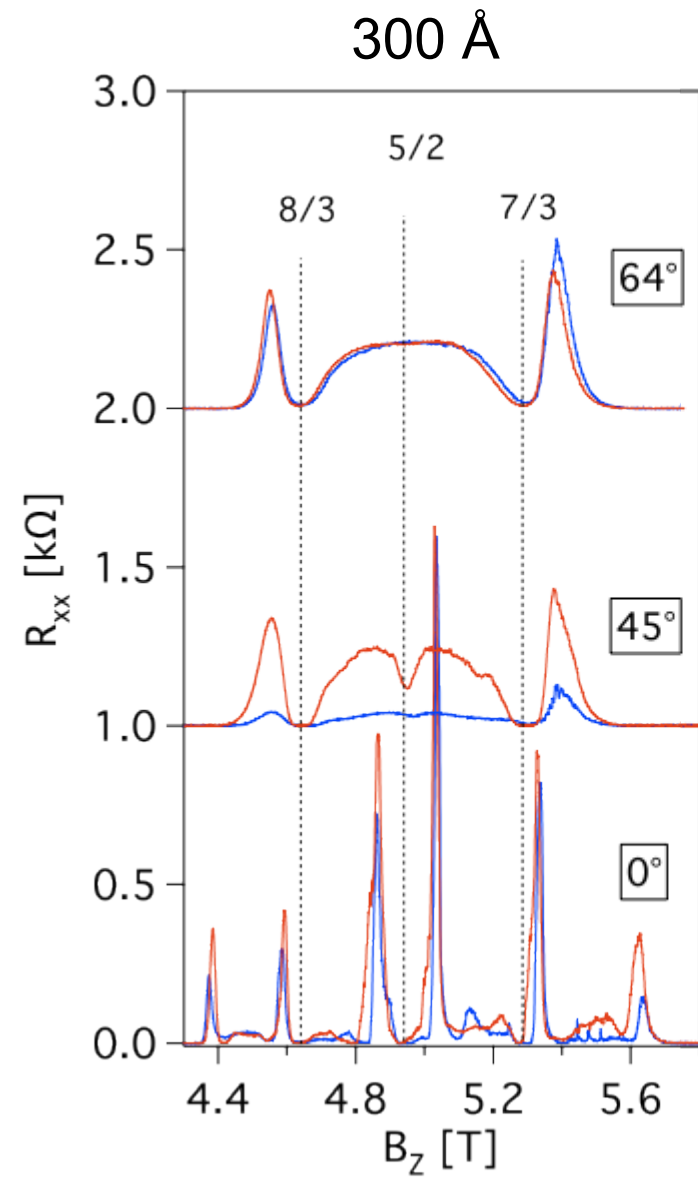
$$n = 3.0 \times 10^{11} \text{ cm}^{-2}$$
$$\mu = 30 \text{ M}$$

*Higher density: Again the nearest empty level is $N=1$ LL in **second** subband*

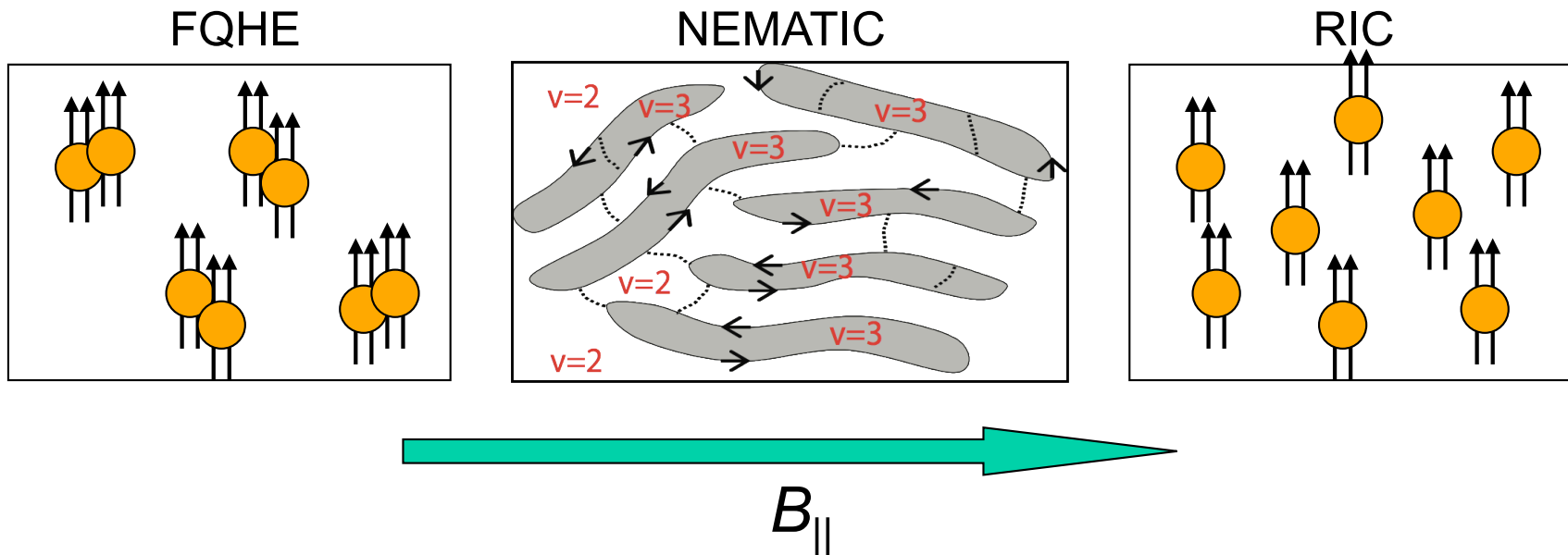
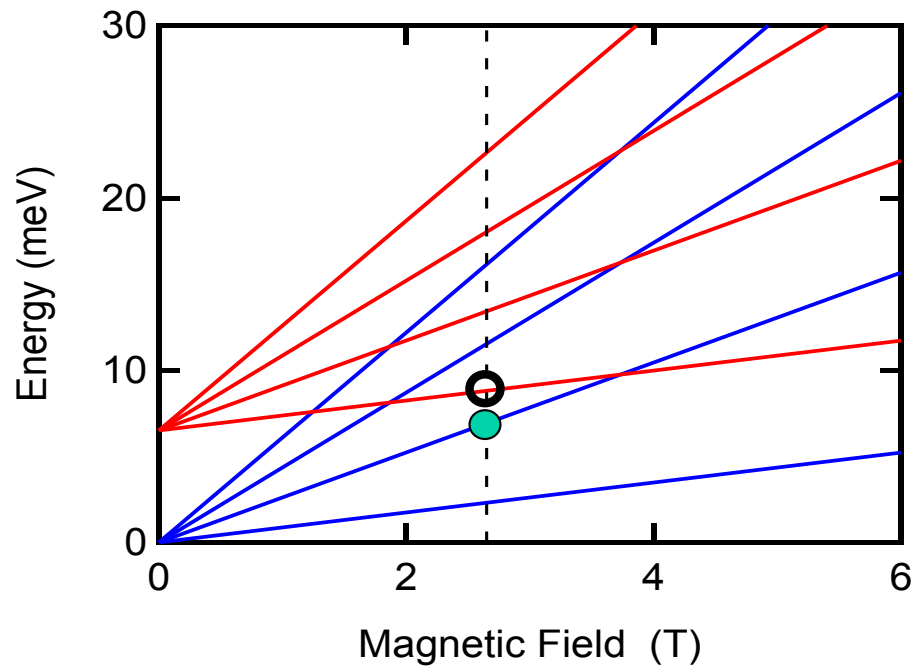
The RIC returns!



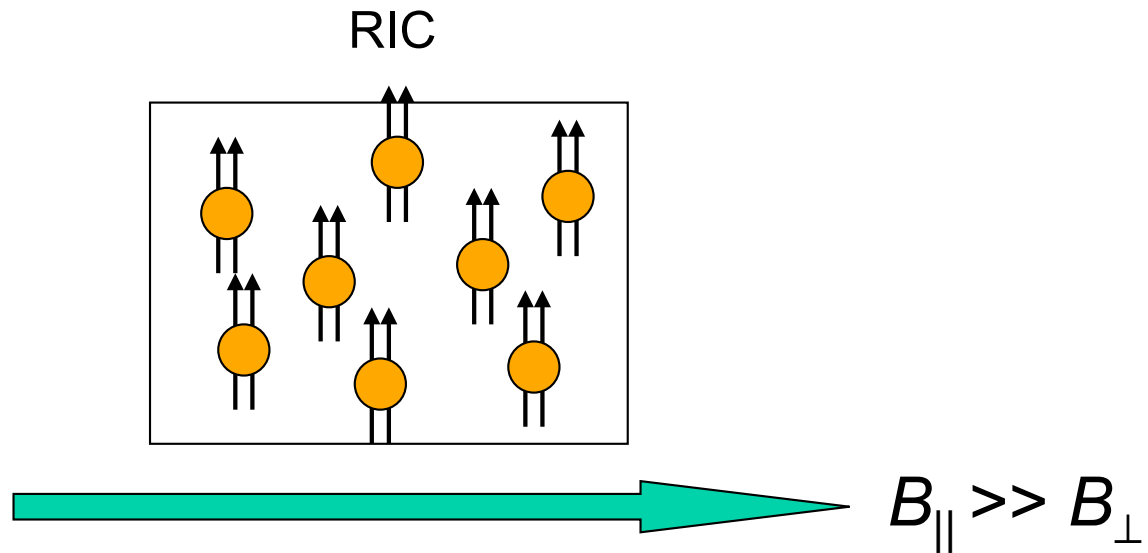
$B_{||}$



Subband-LL alignment is the key to RIC at 5/2

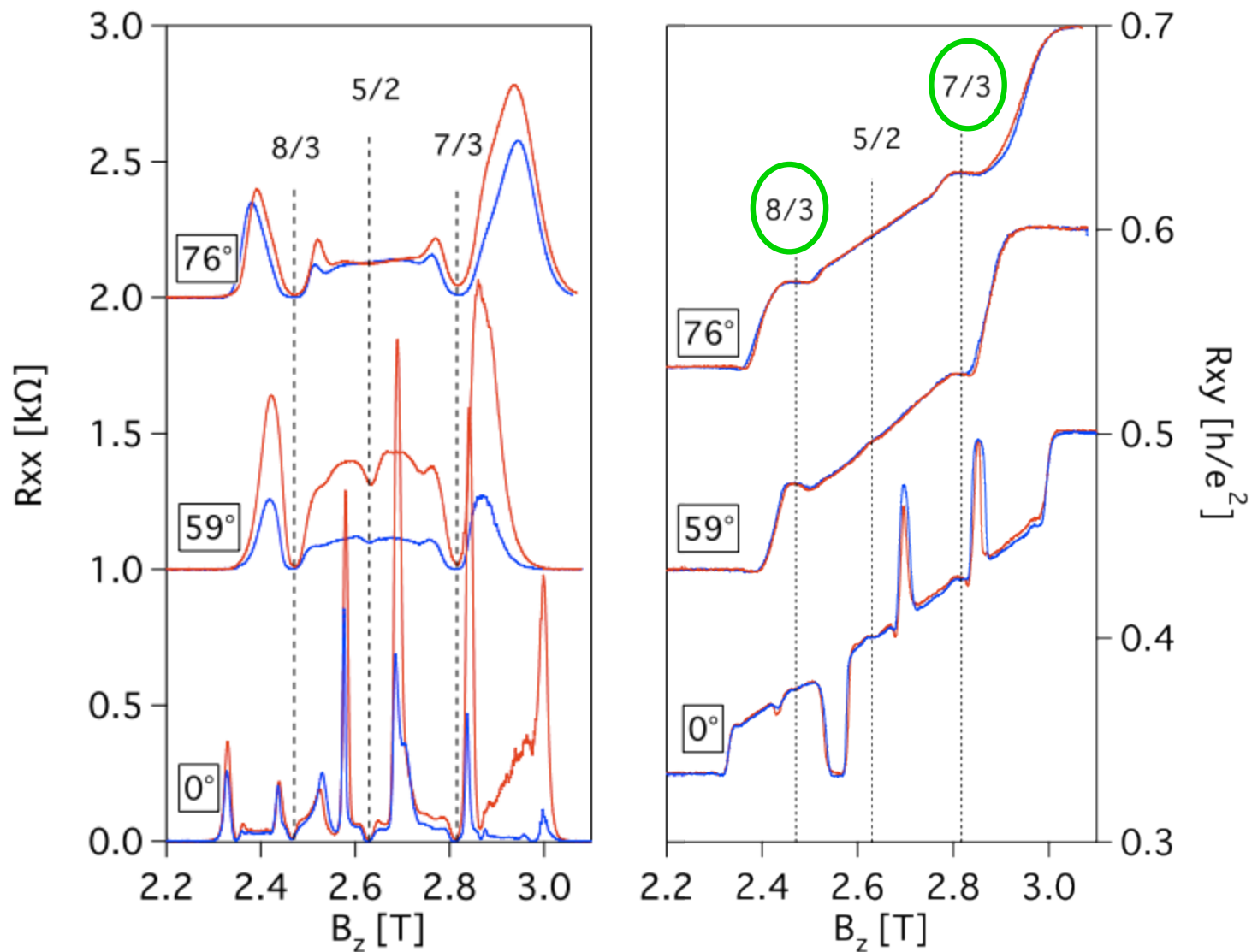


What is the nature of the RIC phase?

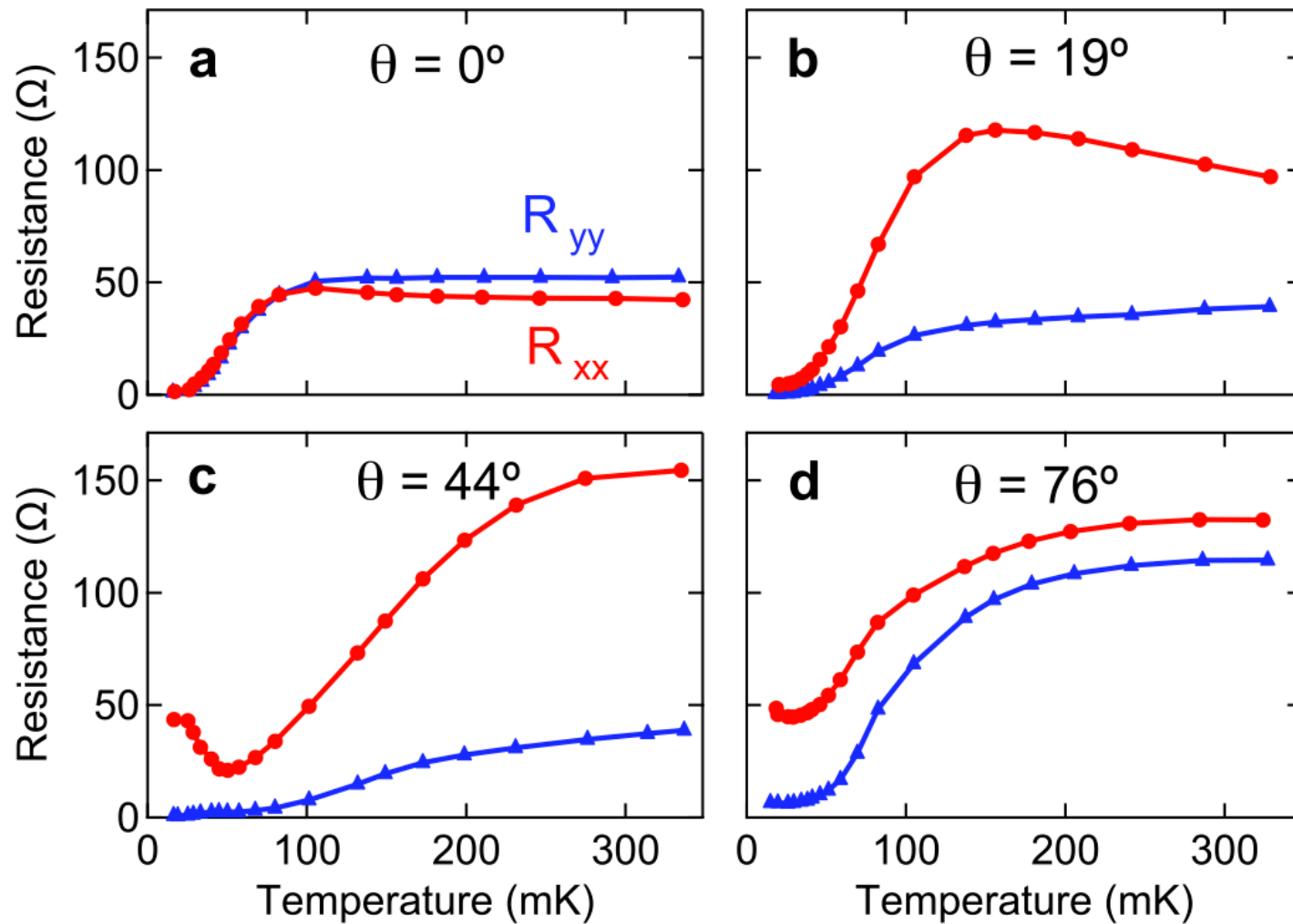


- *Is it a CF Fermi liquid?*
- *If so, is the Fermi surface elliptical?*
- *Why is the transport so isotropic?*

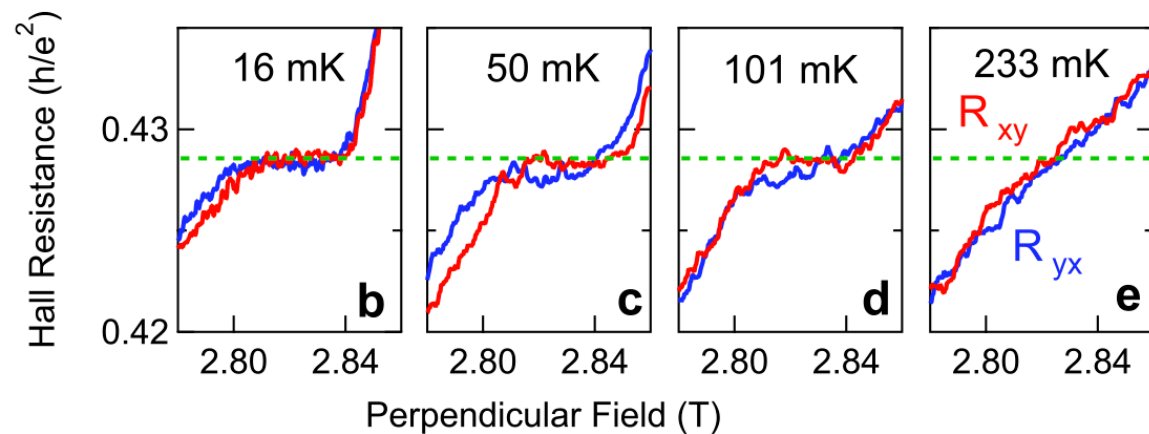
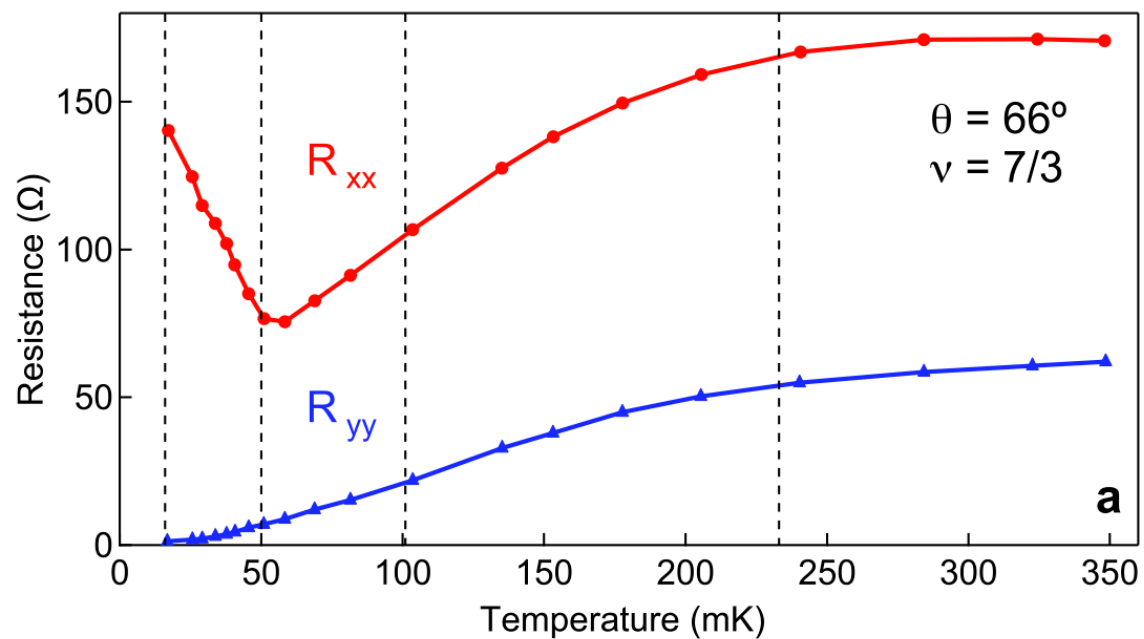
Robust $7/3$ & $8/3$ FQHE at high tilt in 400Å sample



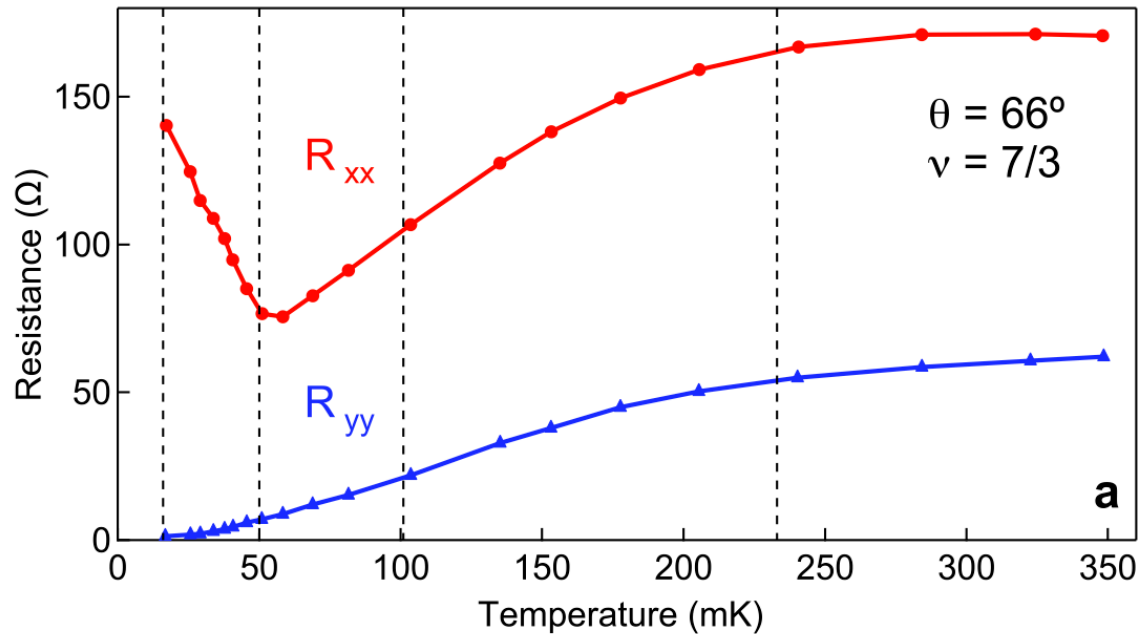
Anisotropic resistivity at 7/3



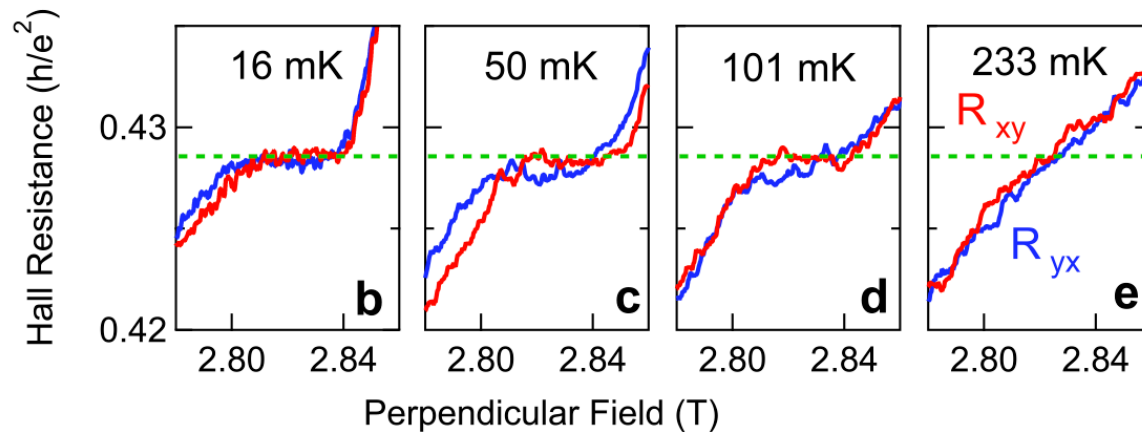
Competing orders at $\nu = 7/3$ in 40 nm sample?



Competing orders at $\nu = 7/3$ in 40 nm sample?



Musaelian & Joynt,
J. Phys. Cond. Matt. 8, L105 (1996).

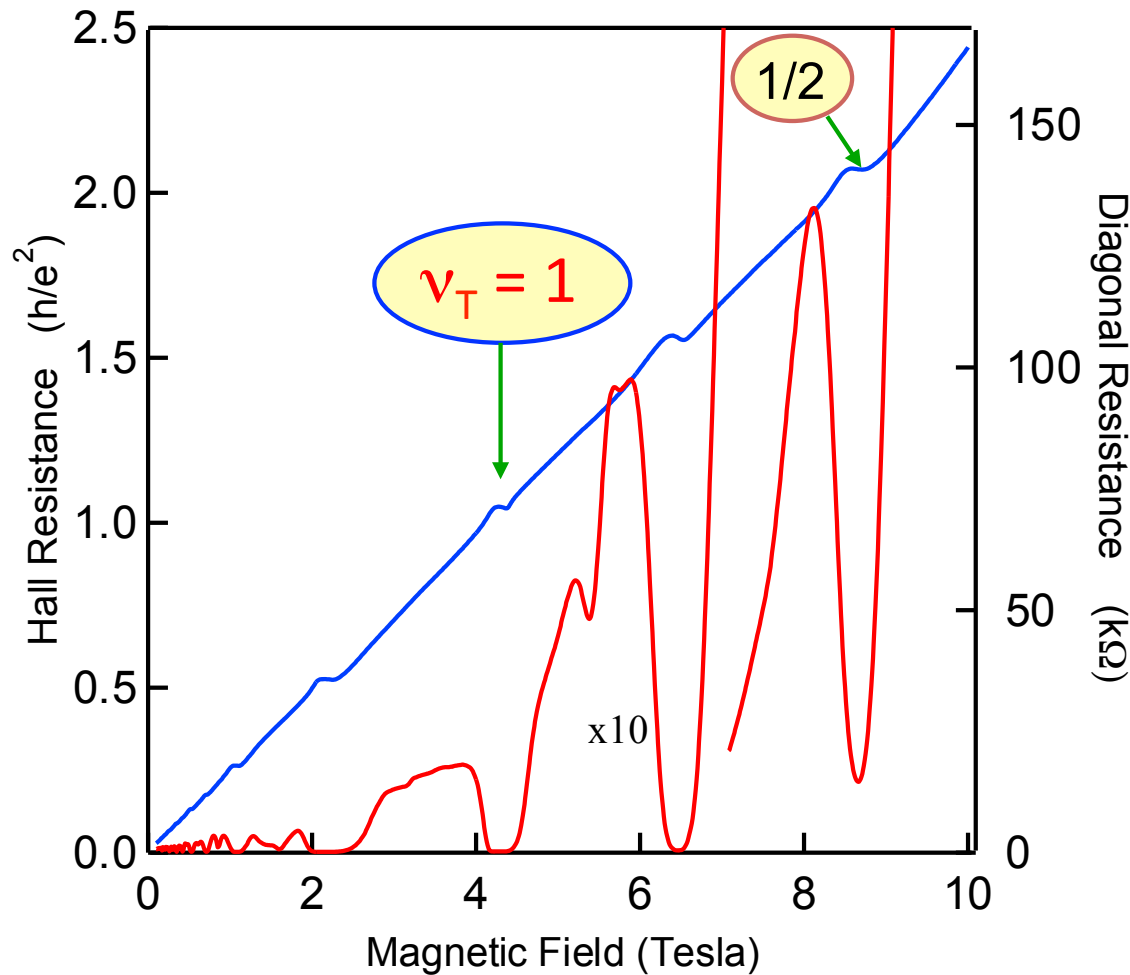


Mulligan, Nayak, and Kachru
PRB 82, 085102 (2010).

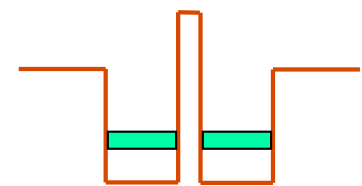
*Hall plateau with stripe-like temperature dependence:
A quantum Hall nematic phase?*

Double layer FQHE systems

QHE in double layer 2D systems

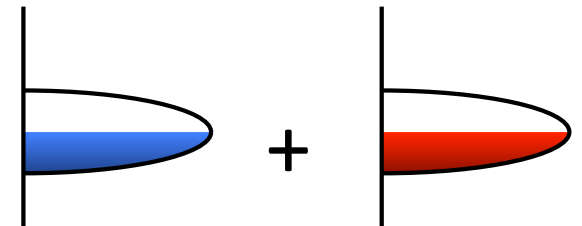


$$\nu_T = \nu_1 + \nu_2$$

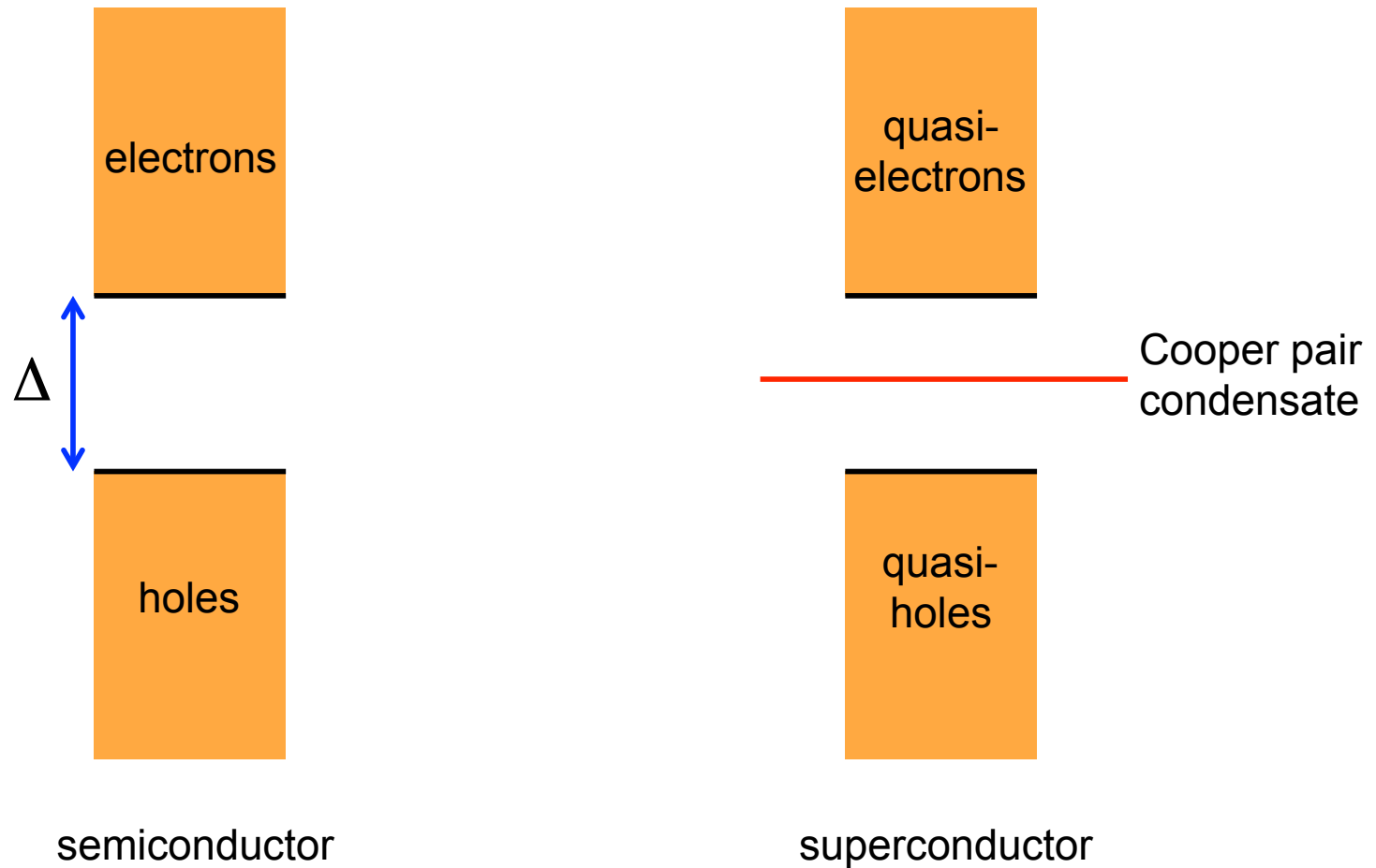


$$\nu_T = 1/2 = 1/4 + 1/4$$

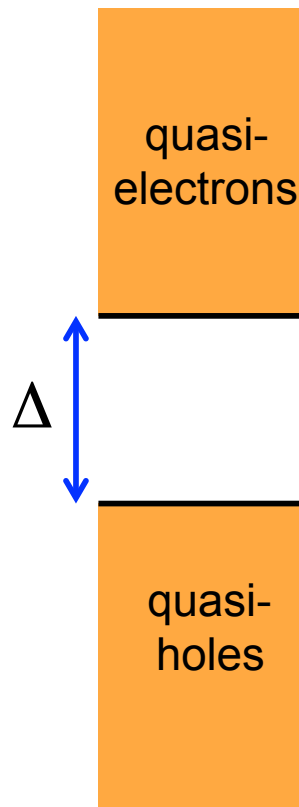
$$\nu_T = 1 = 1/2 + 1/2$$



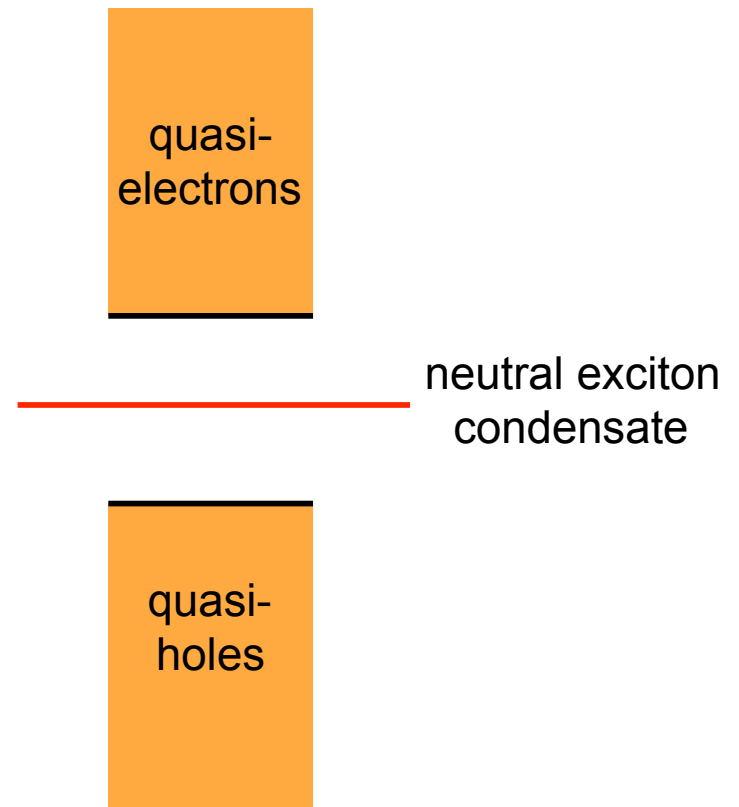
What's special about $\nu_T = 1$ bilayers?



What's special about $\nu_T = 1$ bilayers?



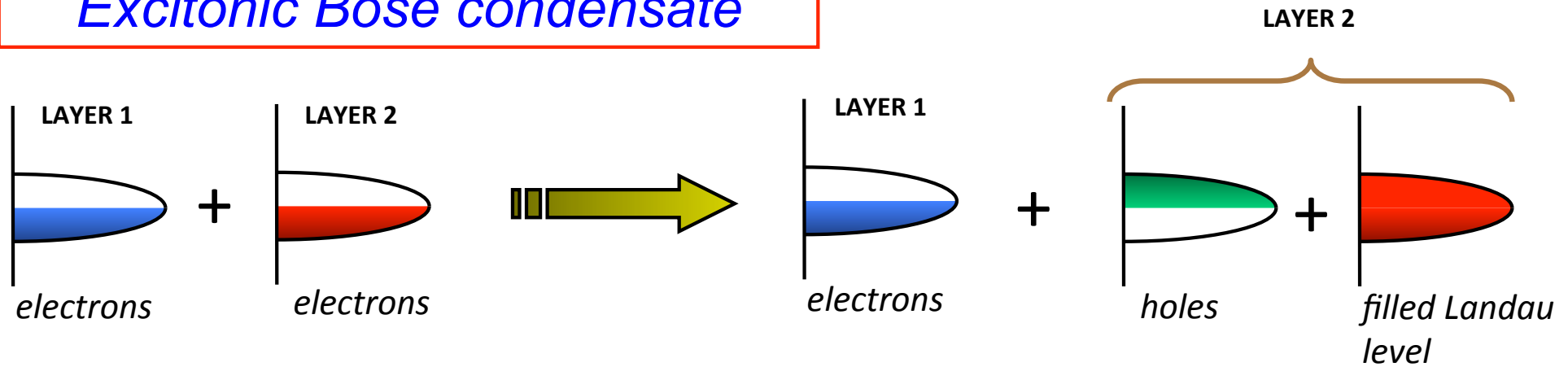
ordinary
QHE



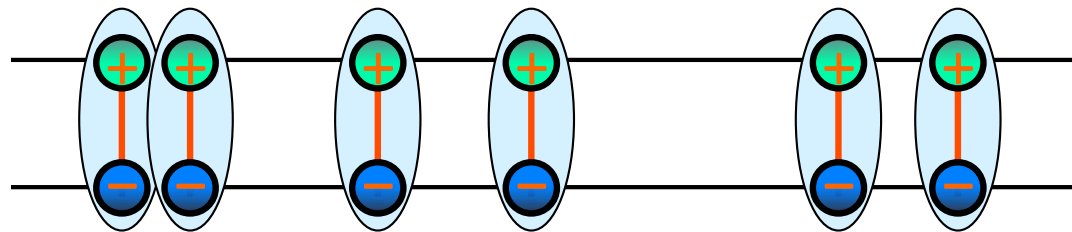
$\nu_T = 1$ bilayer
QHE



Excitonic Bose condensate



$$|\Psi\rangle = \prod_k \frac{1}{\sqrt{2}} \left[1 + \underbrace{e^{i\phi} c_{k,1}^\dagger c_{k,2}}_{\text{exciton creation operator}} \right] |vac'\rangle$$



$$\nabla\phi$$

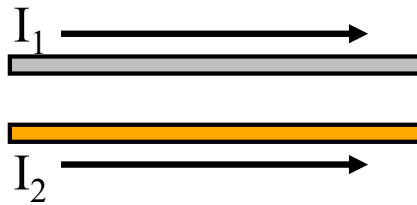


excitonic supercurrents

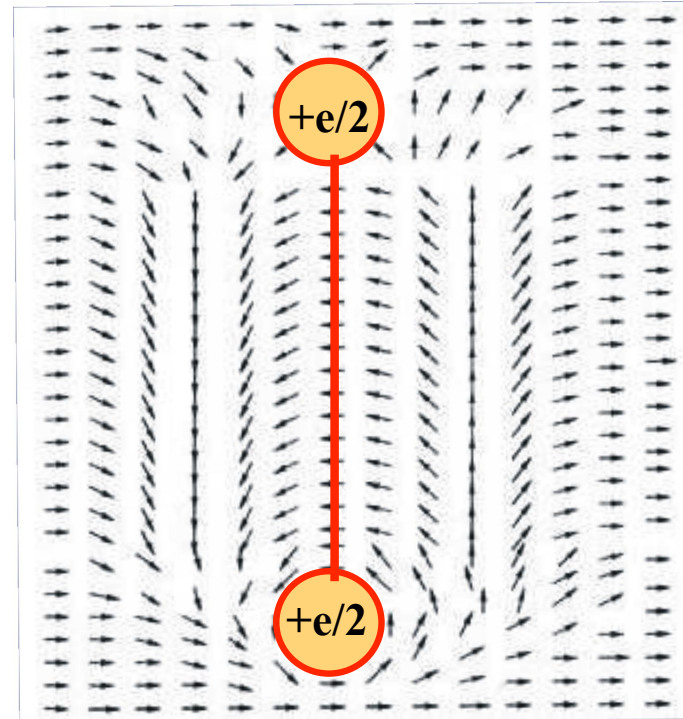


Two transport channels

1. Parallel Transport



meron / anti-meron pair

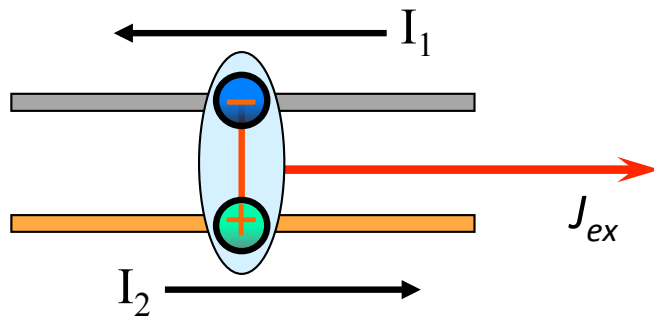


quantized Hall effect



Two transport channels

2. Counterflow Transport



$\nabla\phi = \text{constant}$

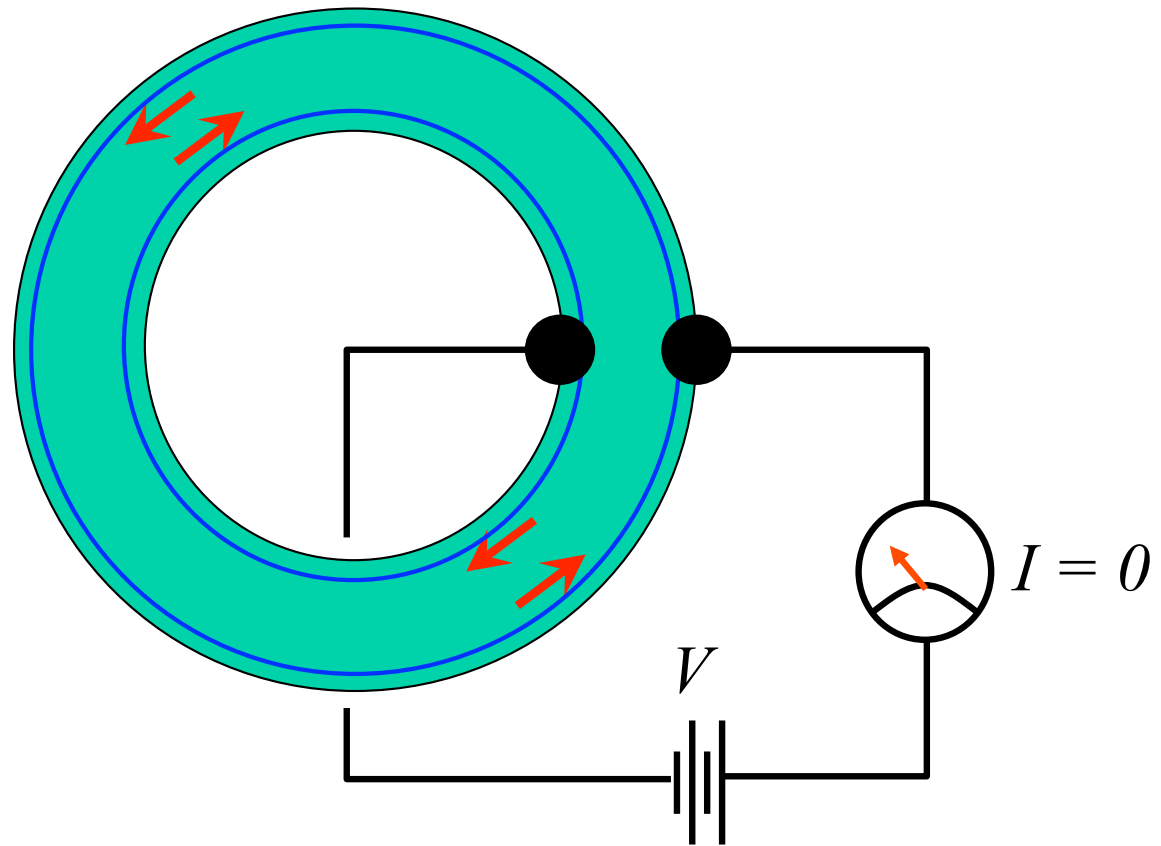
A vector field diagram showing a grid of arrows. The arrows are arranged in a pattern that suggests a constant potential gradient, with some arrows pointing up, some down, and some at an angle, indicating a non-uniform field. Below the vector field, a red arrow points to the right, with the equation $J_{ex} = \rho_s \nabla\phi$ written below it.

$J_{ex} = \rho_s \nabla\phi$

collective exciton transport in condensate



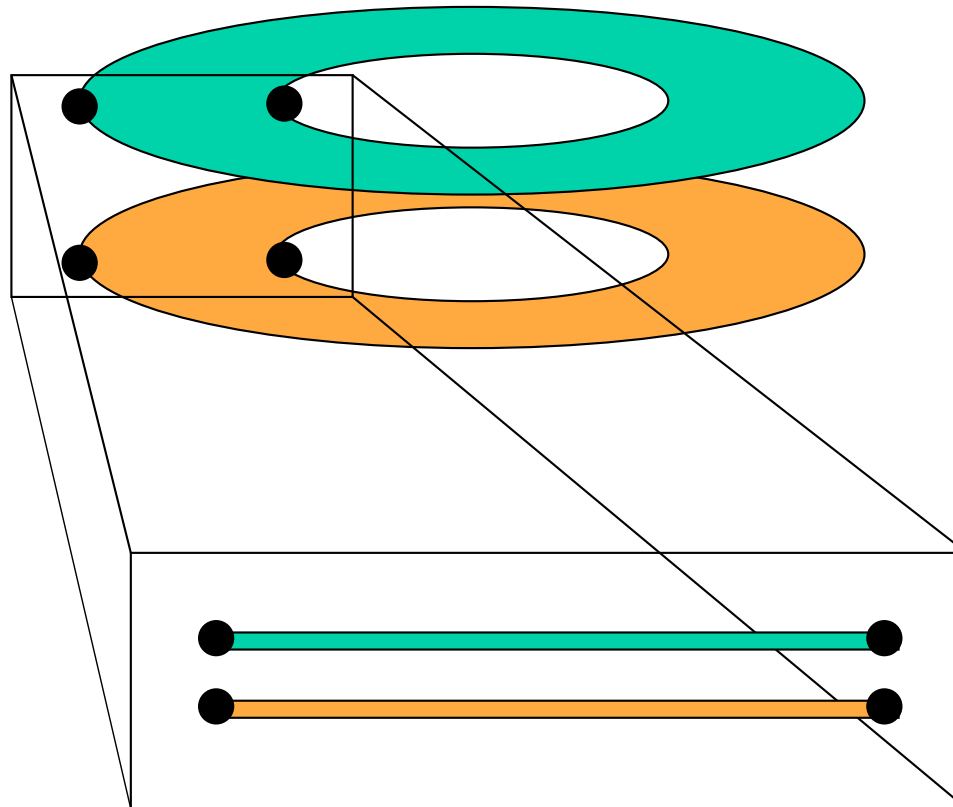
Quantum Hall systems are topological insulators



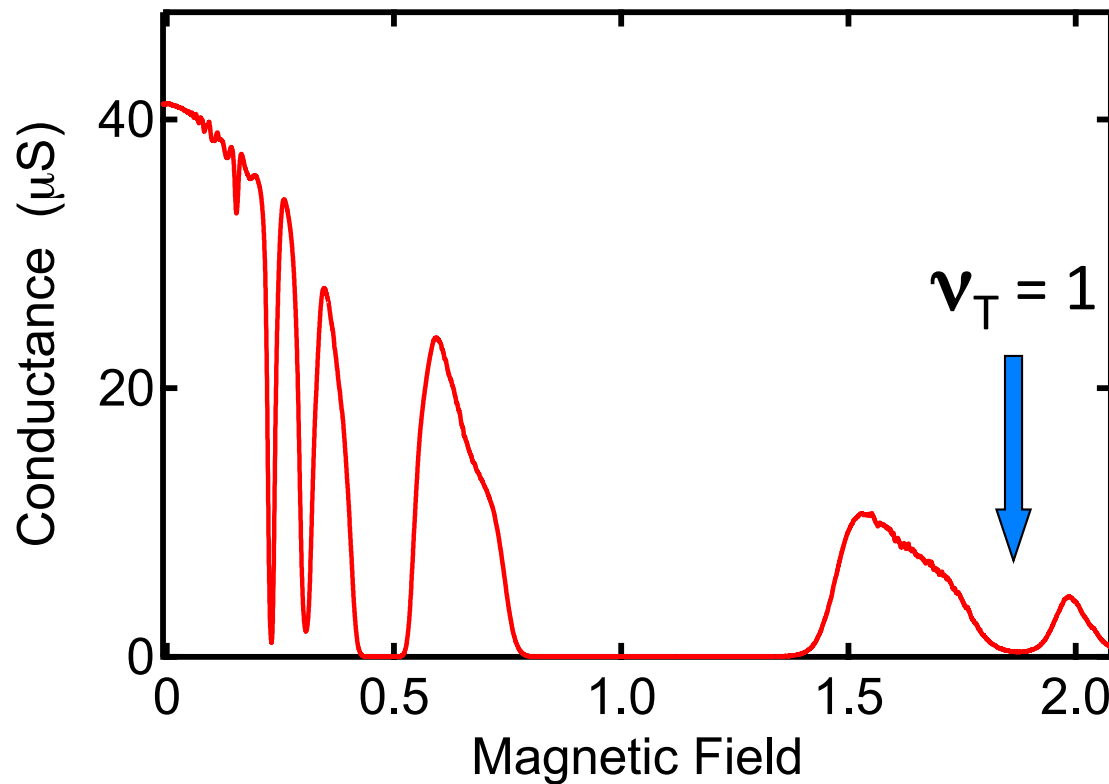
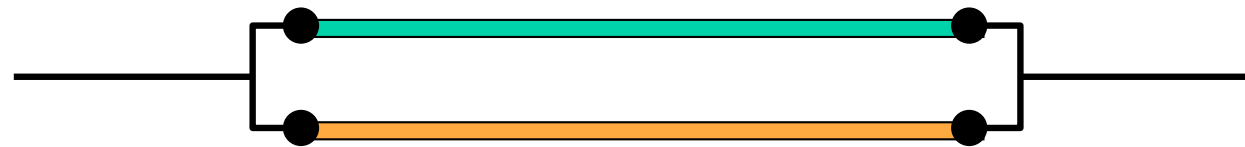
Bulk is opaque to charge transport.



Bilayer Corbino geometry



QHE suppresses parallel charge transport across the bulk



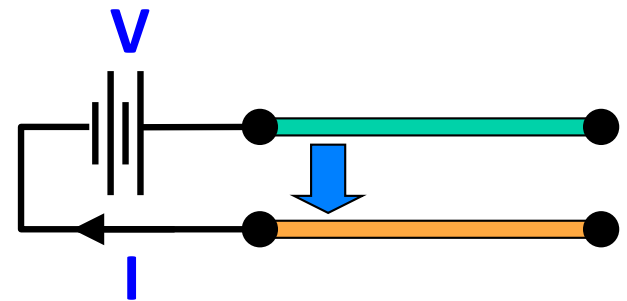
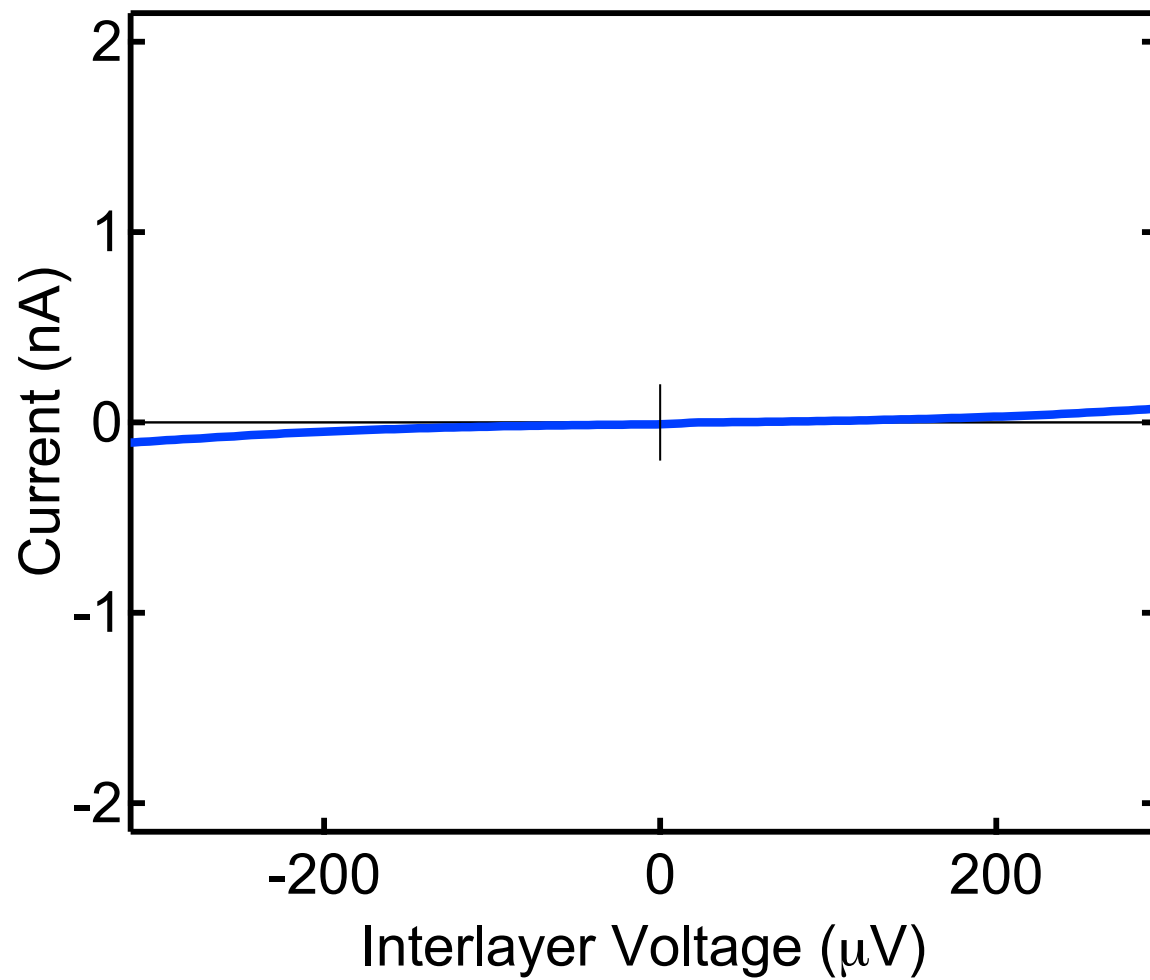
$T = 25 \text{ mK}$

$d/\ell = 1.5$

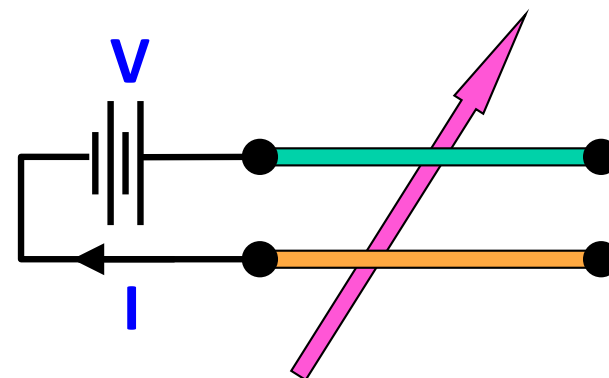
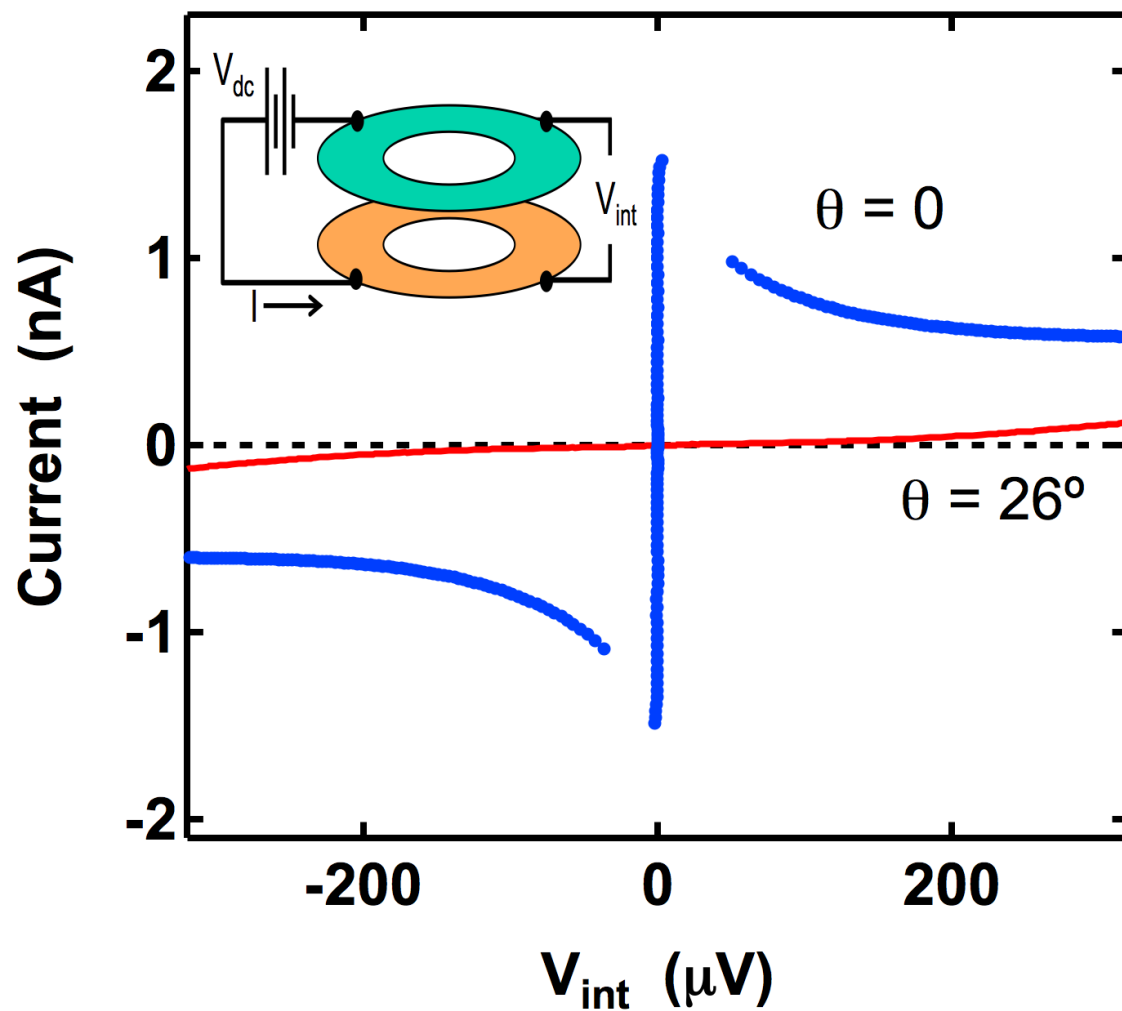
No surprise here.



Tunneling configuration



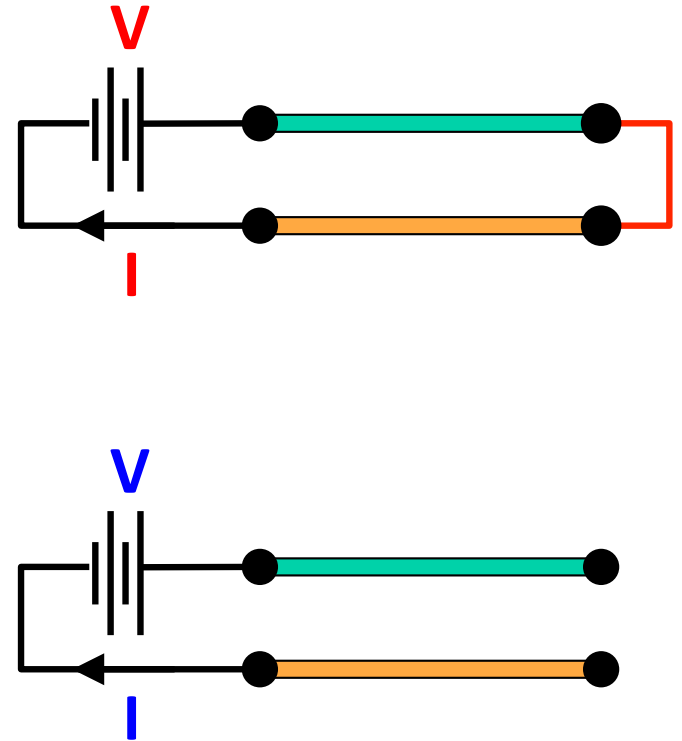
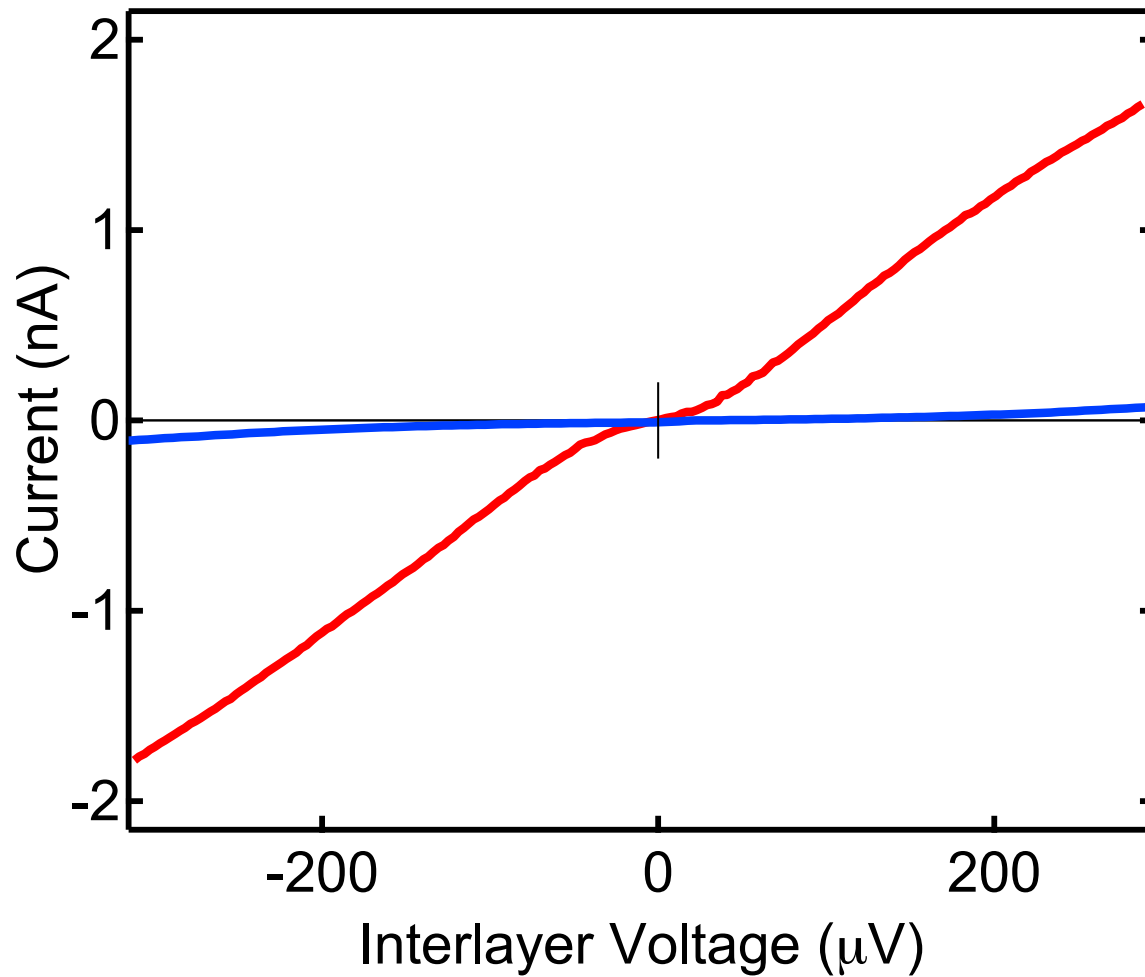
Tunneling configuration



Tunneling intentionally suppressed by tilting.



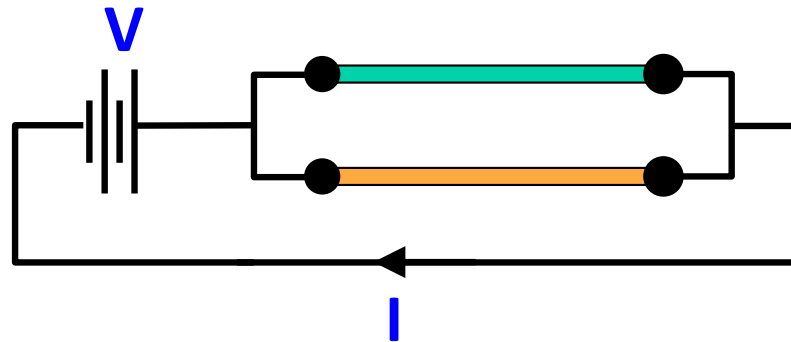
Tunneling vs. Corbino counterflow



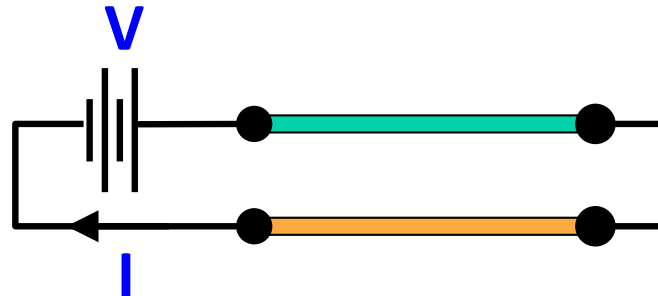
Remote short vastly enhances current.



A paradox?



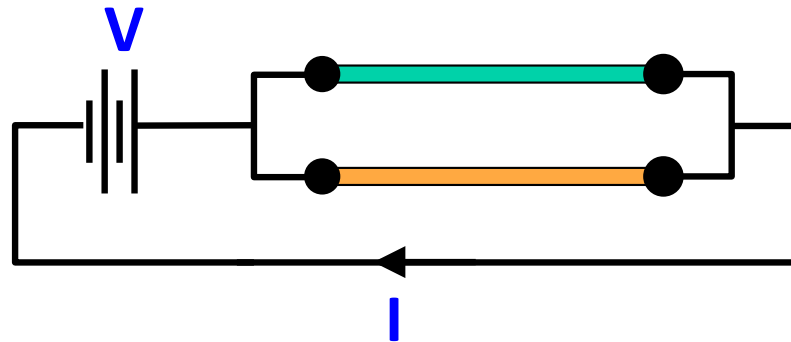
No current



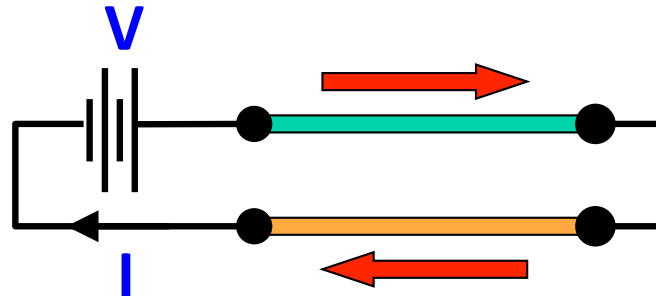
Lots of current



A paradox?



No current



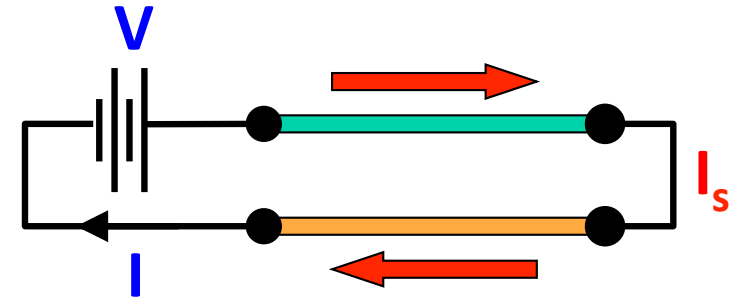
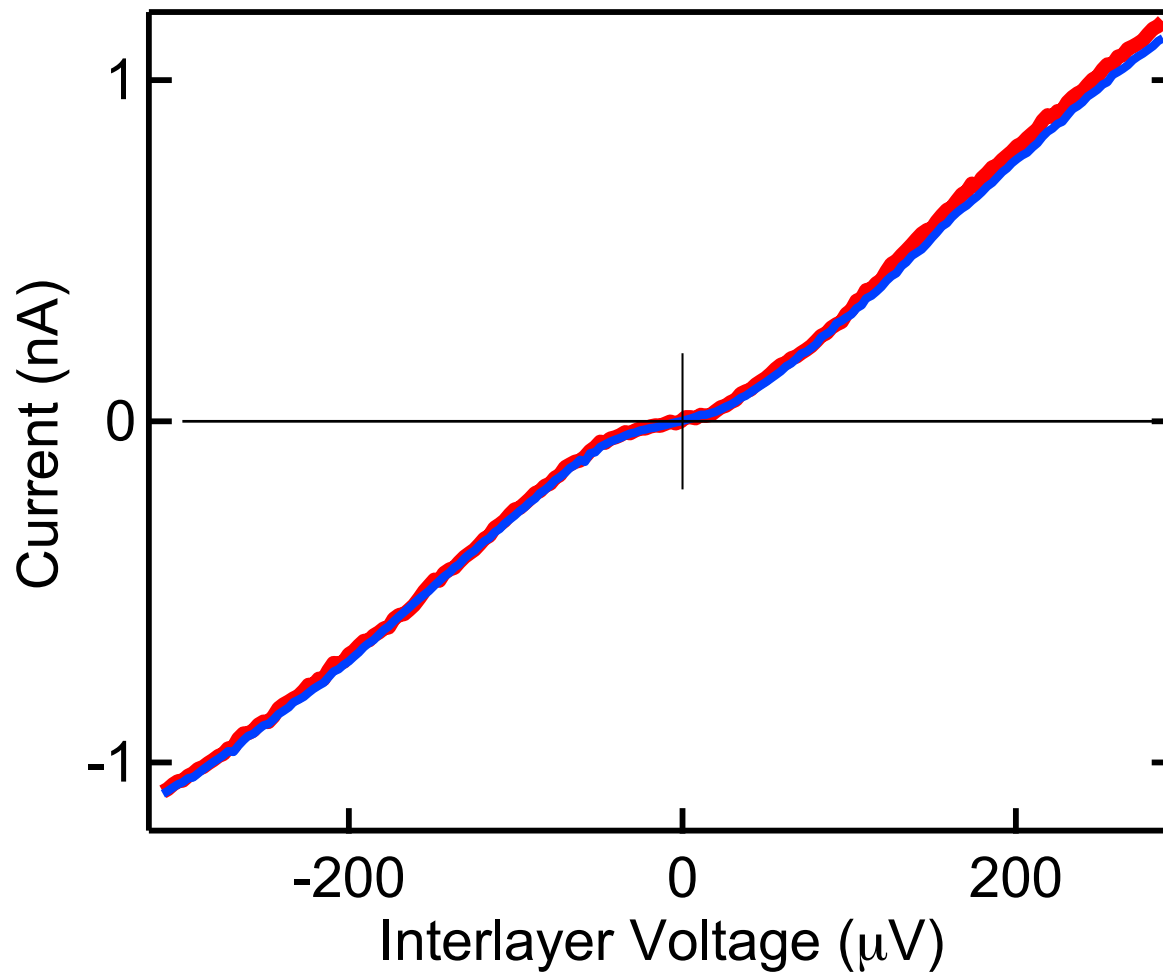
Lots of current

Measure current
here!

Counterflow?



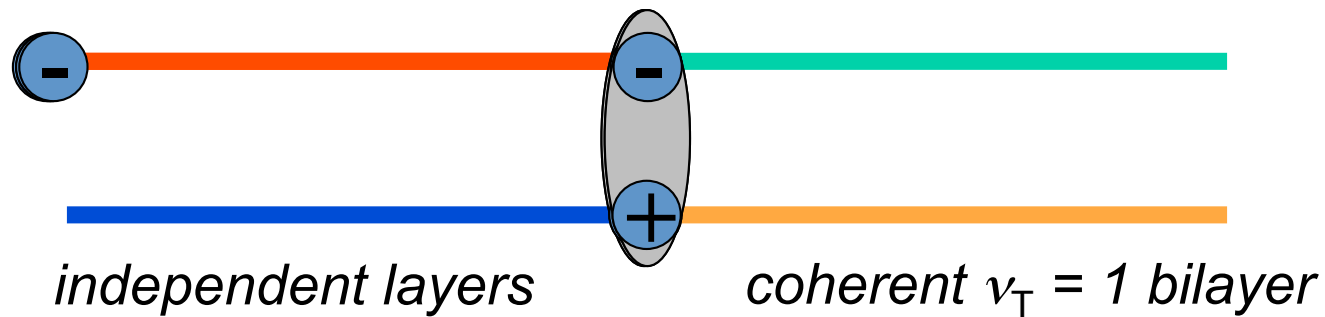
Measuring the shunt current



It IS counterflow.



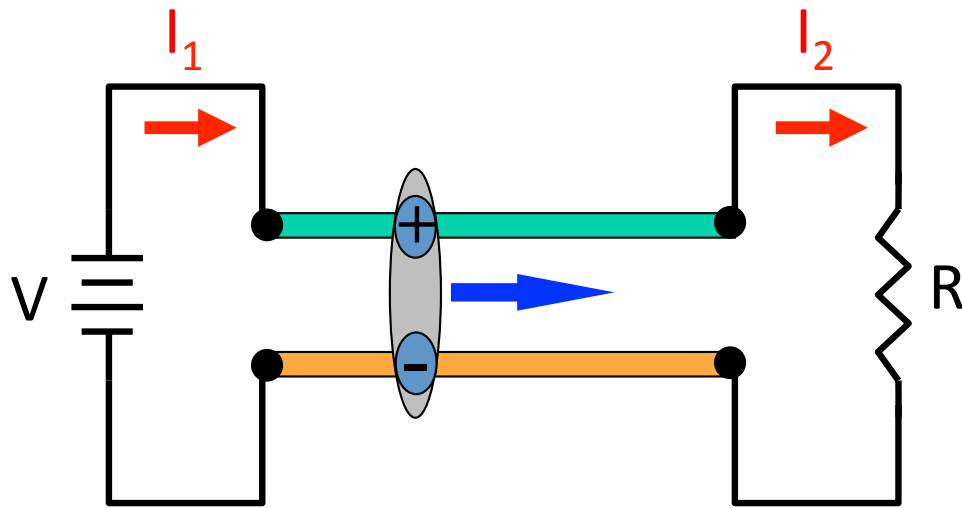
The mechanism



*Excitons are launched and absorbed via Andreev reflection.
Excitons transport energy but not charge.*



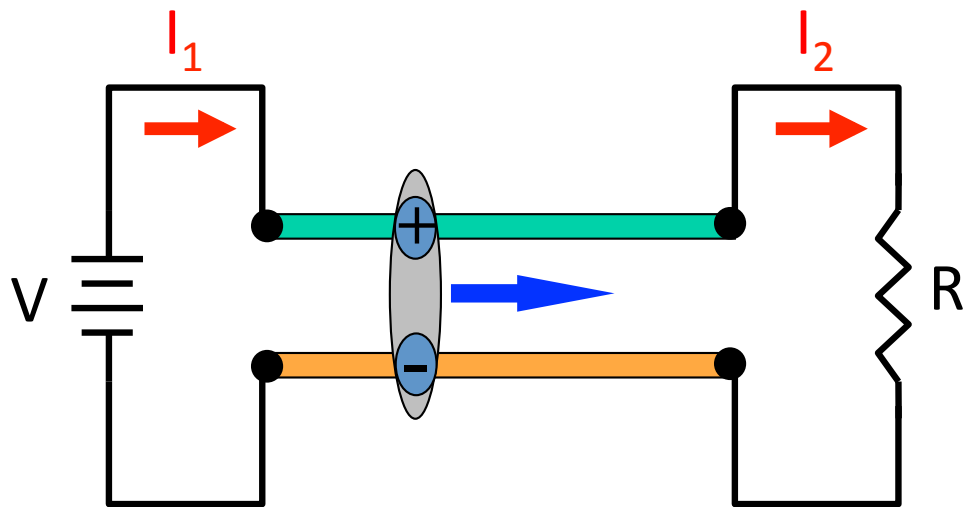
Inducing exciton transport



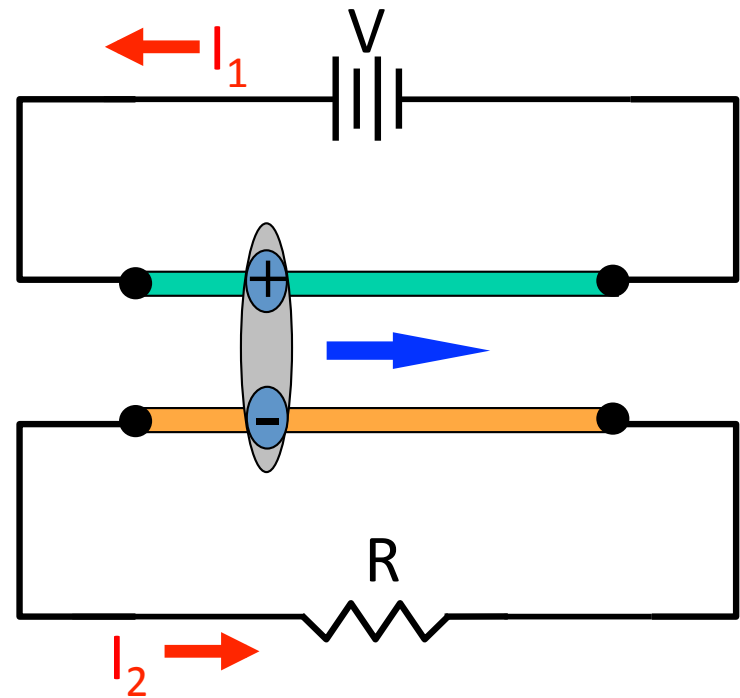
Ideally, $I_1 = I_2$



Inducing exciton transport

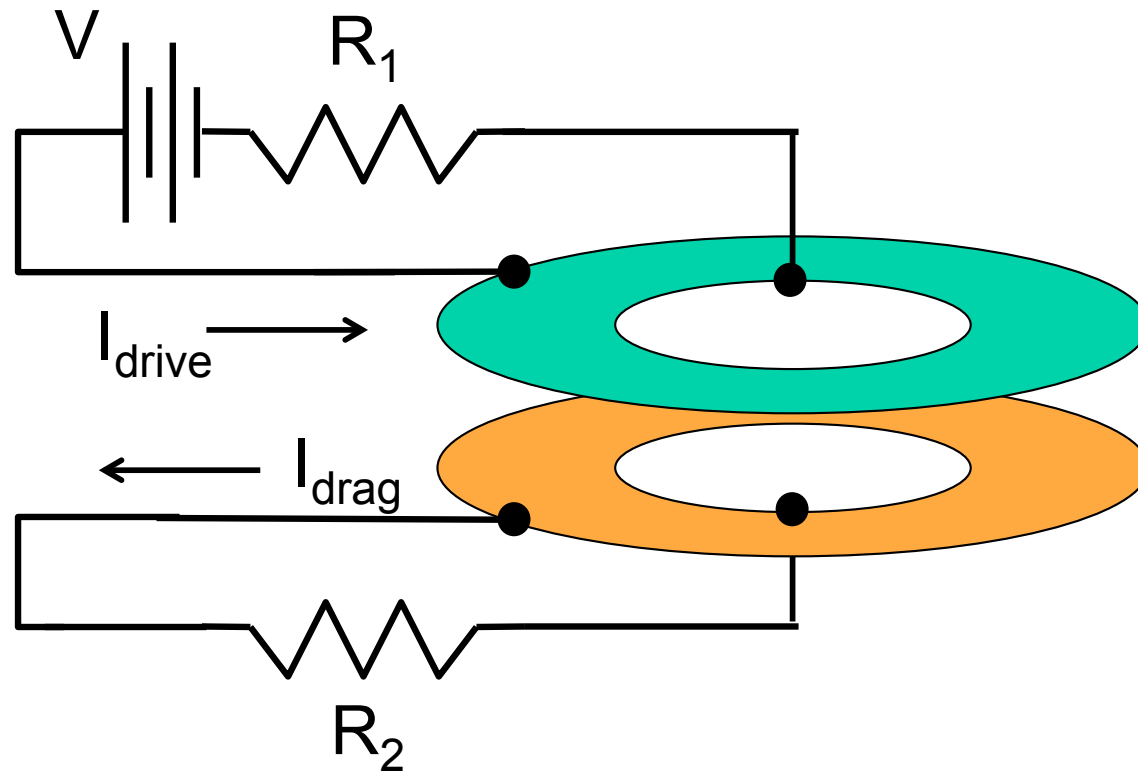


Ideally, $I_1 = I_2$



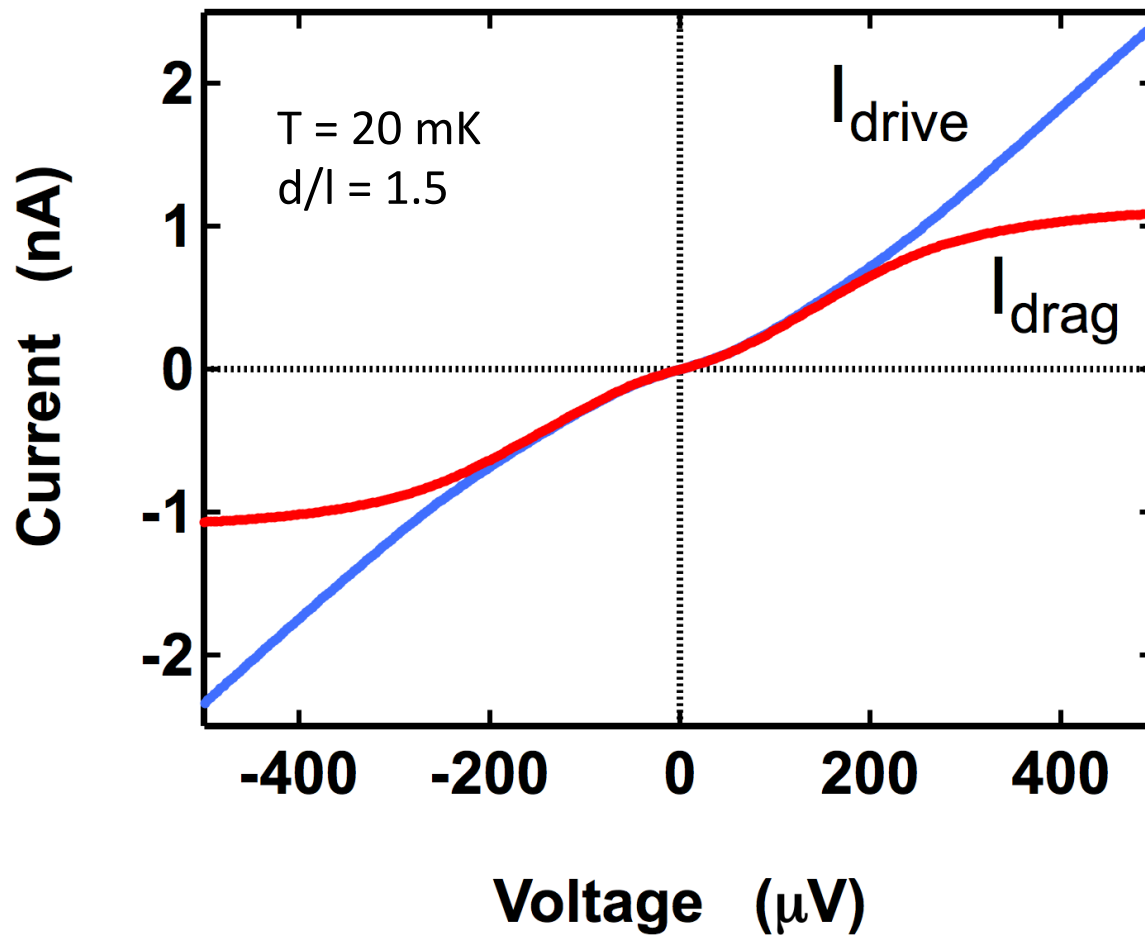
Su & MacDonald 2008

Corbino Coulomb drag



cf. Tiemann et al. PRB 2008

Perfect Coulomb drag



Drag and drive currents equal at small V .
“Perfect” Coulomb Drag



Conclusions

Single layer 2D systems

- *Competing phases at $\nu = 5/2$ in the $N=1$ Landau level.*
- *Details of quantum confinement can be qualitatively important.*
- *Evidence for an anisotropic FQHE state.*

Double layer 2D systems

- *Direct observation of exciton transport across insulating bulk of the bilayer $\nu_T = 1$ QHE state.*
- *Energy transport without charge transport.*
- *“Perfect” Coulomb drag at low T , d/l , and V .*

